

## A CAUCHY LIKE PROBLEM IN PLANE ELASTICITY\*

DANG DINH ANG, NGUYEN DUNG

*Institute of Applied Mechanics, VAST, Hochiminh City*

**Abstract.** Let  $\Omega$  be a bounded domain in the plane, representing an elastic body. Let  $\Gamma_0$  be a portion of the boundary  $\Gamma$  of  $\Omega$ ,  $\Gamma_0$  being assumed to be parallel to the  $x$  - axis. It is proposed to determine the stress field in  $\Omega$  from the displacements and surface stresses given on  $\Gamma_0$ . Under the assumption of plane stress, it is shown that  $\sigma_x + \sigma_y$  is a harmonic function. An Airy stress function is introduced, from which the stress field is computed.

Consider an elastic body represented by a bounded domain  $\Omega$  in the plane. Let  $\Gamma_0$  be a portion of the boundary  $\Gamma$  of  $\Omega$  assumed to be parallel to the  $x$  - axis (cf. Fig. 1). We propose to determine the stress field in  $\Omega$  from the displacements and surface stresses given on  $\Gamma_0$ .

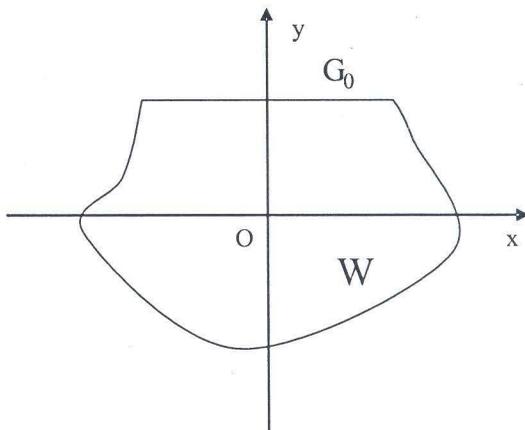


Fig. 1

Cauchy like problems in plane elasticity are treated in [1], [4] and others (cf. References). For a derivation of basic relations on stresses and displacements, we follow {TG}. Assume plane stress. We denote the displacements in the  $x$  - and  $y$  -directions respectively

\*Supported by the Council for Natural Sciences of Vietnam.

by  $u$  and  $v$  and the stress components by  $\sigma_x$ ,  $\sigma_y$  and  $\tau_{xy}$ . Now, we have

$$\varepsilon_x = \frac{\partial u}{\partial x}; \quad \varepsilon_y = \frac{\partial v}{\partial y}; \quad \gamma_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}, \quad (1)$$

$$\begin{cases} \varepsilon_x = \frac{1}{E}(\sigma_x - \nu\sigma_y), & \varepsilon_y = \frac{1}{E}(\sigma_y - \nu\sigma_x), \\ \gamma_{xy} = \frac{1}{G}\tau_{xy} = \frac{2(1+\nu)}{E}\tau_{xy} \end{cases} \quad (2)$$

In the absence of body forces, we have

$$\frac{\partial \sigma_x}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} = 0, \quad (3)$$

$$\frac{\partial \sigma_y}{\partial y} + \frac{\partial \tau_{xy}}{\partial x} = 0. \quad (4)$$

These are the equilibrium equations for our problem.

We now derive the compatibility equations. We have

$$\varepsilon_x = \frac{\partial u}{\partial x}, \quad \varepsilon_y = \frac{\partial v}{\partial y}, \quad \gamma_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \quad (5)$$

from which we get upon differentiating with respect to  $y$ , then with respect to  $x$

$$\frac{\partial^2 \varepsilon_x}{\partial y^2} + \frac{\partial^2 \varepsilon_x}{\partial x^2} = \frac{\partial^2 \gamma_{xy}}{\partial x \partial y}. \quad (6)$$

This differential relation, called the condition of compatibility, must be satisfied by the strain components. By using Hooke's law, the condition (6) can be transformed into a relation between the components of stress.

We have

$$\begin{cases} \varepsilon_x = \frac{1}{E}(\sigma_x - \nu\sigma_y), & \varepsilon_y = \frac{1}{E}(\sigma_y - \nu\sigma_x), \\ \gamma_{xy} = \frac{1}{G}\tau_{xy} = \frac{2(1+\nu)}{E}\tau_{xy}. \end{cases} \quad (7)$$

Substituting into (6), we find

$$\frac{\partial^2}{\partial y^2}(\sigma_x - \nu\sigma_y) + \frac{\partial^2}{\partial x^2}(\sigma_y - \nu\sigma_x) = 2(1+\nu) \frac{\partial^2 \tau_{xy}}{\partial x \partial y}. \quad (8)$$

Differentiating equation (3) with respect to  $x$ , equation (4) with respect to  $y$  and adding together, we find

$$\frac{\partial^2 \tau_{xy}}{\partial x \partial y} = -\frac{\partial^2 \sigma_x}{\partial x^2} - \frac{\partial^2 \sigma_y}{\partial y^2}. \quad (9)$$

Substituting into (8), the compatibility equation in terms of stress components becomes

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) (\sigma_x + \sigma_y) = 0. \quad (10)$$

Before going further, we make the assumption that there exists an open set  $\Omega_0$  with  $\Gamma_0 \subset \Omega_0 \cap \Omega$  such that the stress field is analytically continued to  $\Gamma_0$ . Let

$$\begin{aligned} l\sigma_x + m\tau_{xy} &= \bar{X} \\ m\sigma_y + l\tau_{xy} &= \bar{Y} \end{aligned} \quad \text{on } \Gamma_0, \quad (11)$$

where  $l$ ,  $m$  are the  $x$  and  $y$  - components of the exterior normal to  $\Gamma_0$  and furthermore

$$\begin{aligned} u &= \bar{u}(x) \\ v &= \bar{v}(x) \end{aligned} \quad \text{on } \Gamma_0. \quad (12)$$

It can be shown that

$$\sigma_x \quad \text{and} \quad \frac{\partial \sigma_x}{\partial y} \quad \text{are known on } \Gamma_0, \quad (13)$$

$$\sigma_y \quad \text{and} \quad \frac{\partial \sigma_y}{\partial y} \quad \text{are known on } \Gamma_0. \quad (14)$$

Thus  $\sigma_x + \sigma_y$  is seen as solution of a Cauchy problem on  $\Omega$ . As is well-known, the problem admits at most one solution. It is also known that the problem is ill-posed. Since by (10)  $\sigma_x + \sigma_y$  is harmonic on  $\Omega$ , it is analytic on  $\Gamma_0$ . Thus if  $z_n = (x_n, k)$   $n = 1, 2, \dots$  is a sequence of points of  $\Gamma_0$  with  $x_i \neq x_j$  for  $i \neq j$  and accumulating at a point interior to  $\Gamma_0$ , then  $\sigma_x + \sigma_y$  is uniquely determined by its values on  $(z_n)$ . Hence the Cauchy problem for the Laplace equation can be formulated as a moment problem, it has been regularized by various methods (cf. e.g. [2] chapter 6).

Now we introduce the Airy stress function as follows. Let  $f = \sigma_x + \sigma_y$ , we define

$$\varphi(x, y) = \frac{1}{4\pi} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} f(\xi, \eta) \ln[(x - \xi)^2 + (y - \eta)^2] d\xi d\eta, \quad (15)$$

where  $f$  is set equal to 0 in the complement of  $\Omega$ . We let

$$\bar{\sigma}_x = \frac{\partial^2 \varphi}{\partial y^2}, \quad \bar{\sigma}_y = \frac{\partial^2 \varphi}{\partial x^2}, \quad \bar{\tau}_{xy} = \frac{\partial^2 \varphi}{\partial x \partial y}. \quad (16)$$

It can be checked that  $\bar{\sigma}_x$ ,  $\bar{\sigma}_y$  and  $\bar{\tau}_{xy}$  satisfy the equilibrium equations and the compatibility relation. Hence

$$\bar{\sigma}_x = \sigma_x, \quad \bar{\sigma}_y = \sigma_y, \quad \bar{\tau}_{xy} = \tau_{xy}. \quad (17)$$

## REFERENCES

1. D. D. Ang, D. D. Trong and Yamamoto, A Cauchy Like Problem in Plane Elasticity: Regularization by Quasi-reversibility with Error Estimates, *Viet J. Math.* **32** (2) (2004) 197-208.
2. D. D. Ang, R. Gorenflo, V. K. Le and D. D. Trong, *Moment Theory and Some Inverse Problem in Potential Theory and Heat Conduction*, Lecture Notes in Mathematics, Springer Berlin 2002.
3. S. P. Timoshenko and J. N. Goodier, *Theory of Elasticity*, Mc Graw-Hill, New York, 1951.
4. D. D. Ang, C. D. Khanh and M. Yamamoto, A Cauchy Like Problems in Plane Elasticity: A Moment Theoretic Approach, *Vietnam J. Math.* **32** (2004) 19-32.
5. D. D. Ang, M. Ikeheta, D. D. Trong and M. Yamamoto, Unique continuation for a Stationary Isotropic Lamé System with Variable Coefficients, *Comm. Partial Diff. Eqs.* **23** (142) (1998) 371-385.

## BÀI TOÁN TỰA CAUCHY TRONG ĐÀN HỒI PHẪNG

Xét một vật thể đàn hồi biểu diễn bởi một miền  $\Omega$  bị chặn trong mặt phẳng. Gọi  $\Gamma_0$  là một phần của biên  $\Gamma$  của miền  $\Omega$ ,  $\Gamma_0$  được giả thiết là song song với trục tọa độ  $x$ . Vấn đề đặt ra là xác định trường ứng suất trong  $\Omega$  từ các chuyển vị và ứng suất mặt cho trước trên  $\Gamma_0$ . Dưới các giả thiết của ứng suất phẳng, bài báo đã chỉ ra  $\sigma_x + \sigma_y$  là một hàm điều hòa. Một hàm ứng suất Airy được giới thiệu để từ đó có thể tính được trường ứng suất .