MONTE CARLO SIMULATION OF VAPOR - LIQUID EQUILIBRIA OF HYDROGEN USING AB INITIO INTERMOLECULAR POTENTIALS

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ABSTRACT

The vapor-liquid equilibria of pure fluid hydrogen were predicted by Gibbs ensemble Monte Carlo simulation techniques using our two different ab initio intermolecular pair potentials. The ab initio intermolecular interaction potentials were obtained from coupled-cluster calculations, using the CCSD(T) level of theory and Dunning's correlation consistent basis sets aug-cc-pVmZ (m = 2, 3) [9]. The phase diagram, critical properties, thermodynamic properties, vapor pressures and orthobaric densities based on them are found to be in very good agreement with experimental data.

Keywords: Vapor-liquid equilibria, Gibbs ensemble Monte Carlo simulation, ab initio potentials.

I - INTRODUCTION

Hydrogen and the mixture hydrogen-oxygen are used in several industrial applications. It could become the most important energy carrier of tomorrow. Liquid hydrogen, oxygen are the usual liquid fuels for rocket engines [1]. The National Aeronautics and Space Administration (NASA) is the largest user of liquid hydrogen in the world [2].

Computer simulations have become indispensable tools for studying pure fluids and fluid mixtures. One of the first attempts Nasrabad and Deiters predicted phase highpressure vapour-liquid phase equilibria of noblegas mixtures [3,4] from the global simulations using the intermolecular potentials. Leonhard and Deiters used a 5-site Morse potential to represent the pair potential of nitrogen [5] and were able to predict vapour pressures and orthobaric densities. Naicker et al. developed the 3-site pair potentials for hydrogen chloride [6]; they predicted successfully the vapourliquid phase equilibria of hydrogen chloride with GEMC (Gibbs Ensemble Monte Carlo simulations [7]).

In this work we report the simulation results of the vapor-liquid equilibria for the pure fluid hydrogen using Gibbs Ensemble Monte Carlo (GEMC) simulation techniques with our new *ab initio* intermolecular pair potentials resulting from quantum mechanical calculations at a sufficiently high level of approximation of dimer H₂-H₂ [8]. The phase equilibrium results density, vapour pressure and enthalpy of vaporization obtained from GEMC simulation are compared with experimental data and with those from literature data.

II - COMPUTATIONAL TECHNIQUE

1. Simulation details

The Gibbs ensemble Monte Carlo (GEMC) simulation techniques were used for calculating the thermodynamic properties of the pure fluid

hydrogen. The GEMC-NPT simulation was used to calculate the density, and the internal energy of the fluid hydrogen to examine the accuracy of the pair potentials. This simulation was investigated on isobars at 1.0 MPa and 5.0 MPa and for temperatures from 26.0 K to 250 K, respectively. GEMC-NVT simulations were performed to obtain coexisting liquid and vapor densities, and vapor pressures. They were in the temperature range 18.0 K to 32.0 K with an increment 2.0 K.

The two our new 5-site pair potential functions for hydrogen resulting from *ab initio* calculations [8] were used for both simulation cases:

$$u_{1} = \sum_{i=1}^{5} \sum_{i=1}^{5} \left[D_{e}^{ij} e^{-\alpha_{ij}r_{ij}} + f_{1a}(r_{ij}) \sum_{n=6,8,10} \frac{C_{n}^{ij}}{r_{ij}^{n}} + f_{2}(r_{ij}) \frac{q_{i}q_{j}}{4\pi\varepsilon_{0}r_{ij}} \right]$$
(1)

$$u_{2} = \sum_{i=1}^{5} \sum_{i=1}^{5} \left[D_{e}^{ij} e^{-\alpha_{ij}r_{ij}} + f_{1b}(r_{ij}) \sum_{n=6,8,10,12} \frac{C_{n}^{ij}}{r_{ij}^{n}} + f_{2}(r_{ij}) \frac{q_{i}q_{j}}{4\pi\varepsilon_{0}r_{ij}} \right]$$
(2)

With: $f_{1a}(r_{ij}) = (1 + e^{-2(\delta_{ij}r_{ij}-2)})^{-15}$, $f_{1b}(r_{ij}) = 1 - e^{-\delta_{ij}r_{ij}} \sum_{k=0}^{10} \frac{(\delta_{ij}r_{ij})^k}{k!}$ and $f_2(r_{ij}) = 1 - e^{-\beta_{ij}r_{ij}}$

Here the r_{ij} site-site distances, the q_i electric charges of sites, and the C_{ij}^n dispersion coefficients; the leading dispersion term is always proportional to r. The two models differ mostly in the choice of the damping functions $f_{1a}(r_{ij})$ and $f_{1b}(r_{ij})$. The optimal adjustable parameters of these functions were estimated by nonlinear least-square fitting to the *ab initio* interaction energy values, as shown in [8].

Total number of particles N = 512 were used in both GEMC simulations with the standard periodic boundary conditions and the minimum image convention. For GEMC-NVT simulation runs the equilibration between two phase required 1-2 x 10^6 cycles. The simulation parameters were set for 50% acceptance ratios for translations and volume fluctuations. All movements were performed randomly with probabilities. The defined accumulative averages of desired quantities were established within 10³ cycles, after initial equilibration had been reached within 5.0 x 10^4 cycles. The simulation data were exported using block averages with 1000 cycles per block. The statistical errors in the simulation runs were estimated by dividing each run into 100 blocks and taking the largest deviation of a block mean from the total mean as error. The simulations were started with equal densities in two phases. The simulation systems were equilibrated for about 1.0 x 10^6 cycles. The cut-off radius r_c was set to 7.5 for hydrogen. Corrections for long-range interactions for the internal energy were computed by the standard relations [10].

2. Structural properties

The structural properties of the fluid hydrogen were studied for the liquid phase at different temperatures with the GEMC-NVT and -NPT simulations, respectively; in both cases the temperature dependence is shown by site-site pair distribution function gr. Because of the 5-site model of dimer H₂-H₂ was constructed with two sites placed on the atoms H, one site in the center of gravity M, and two sites halfways between the atoms and the center N [8]. Consequently, the 5-site pair correlation functions also consisted of the interactions H-H, N-N, M-M, N-M, H-N and H-M for the fluid hydrogen. These were achieved by simulations. The structural properties of fluid were compared with experimental data and with data from literature, if available.

3. Phase coexistence properties

The critical temperature T_c/K , density $\rho_c/g.cm^{-3}$ and volume $V_c/cm^3.mol^{-1}$ of the pure

fluid hydrogen were derived from least-squares fits to the densities of coexisting phase using the relations of the rectilinear diameter law [10]:

$$\frac{\rho_{1} - \rho_{2}}{2} = \rho_{c} + A(T_{c} - T)$$

$$\rho_{1} - \rho_{2} = B(T - T_{c})^{\beta}$$
(3)

where ρ_1 and ρ_v are the coexistence liquid density and vapor density, β is the critical exponent ($\beta = 0.325$). A and B are adjustable constants. The critical pressure P_c / MPa was calculated with the Antoine equation

$$\ln P = A - \frac{B}{T - C} \tag{4}$$

where A, B and C are Antoine constants.

The relation between vapor pressure, heat of vaporization $\Delta_{vap}H$ and temperature is given by the Clausius-Clapeyron equation:

$$\ln \frac{P_{1}}{P_{2}} = -\frac{\Delta_{vap}H}{R} \left(\frac{1}{T_{1}} + \frac{1}{T_{2}}\right)$$
(5)

For the standard state $P^0 = 0.101$ MPa this relation is rewritten as

$$\ln\frac{P}{P^0} = -\frac{\Delta_{vap}H}{R}\frac{1}{T} + \frac{\Delta_{vap}S}{R}$$
(6)

Here T_1 and T_2 are the temperatures at the

pressure P_1 and P_2 . The slope and the intercept of $\ln P$ with respect to 1/T are proportional to $\Delta_{vap}H$ and $\Delta_{vap}S$.

III - RESULTS AND DISCUSSION

1. Structural properties

This section describes the features of the site-site pair distribution functions resulting from two GEMC-NVT and -NPT simulation techniques for the pure fluid hydrogen. The *ab initio* pair potentials Eqs. 1 and 2 of hydrogen, respectively, were used for those simulations.

The temperature dependence of the radial distribution functions at two different pressures 1.0 MPa and 5.0 MPa is obtained from the global simulations GEMC-NPT and GEMC-NVT. Figures 1 and 2 show simulation results for hydrogen. As 5-site potential models Eq.1 and Eq. 2 were used for these simulations, they consist of the interactions of ghost sites N, M on the molecules, so these were also represented here by 5-site pair correlation functions [8]. The height of site-site pair correlation functions decreased with increasing temperature from 26.0 K to 250.0 K, as shown in Table 1. In general the peaks for the interaction of sites including an atomic nucleus were higher than those without a nucleus.

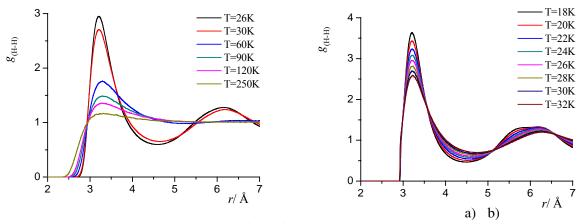


Figure 1: Temperature dependence of g_{H-H} for hydrogen at P = 1.0 MPa a) simulation GEMC-NPT and b) simulation GEMC-NVT, in both cases using pair potential Eq.1

All first peaks of the site-site correlation functions for hydrogen are located between 2.893 Å and 3.205 Å. The second peaks are located between 6.081 Å and 6.234 Å.

Equation	<i>T</i> / K	at 1.0 MPa			at 5.0 MPa				
		g _{H-H}	g _{N-N}	g _{M-M}	g _{H-M}	g _{H-H}	g _{N-N}	g _{M-M}	g _{H-M}
Eq. 1	26.0	2.92	1.97	2.53	2.69	2.79	1.85	2.40	2.56
	30.0	2.69	1.87	2.37	2.51	2.55	1.76	2.24	2.39
	60.0	1.75	1.51	1.67	1.71	1.67	1.44	1.59	1.62
	90.0	1.49	1.35	1.44	1.47	1.42	1.28	1.37	1.40
Eq. 2	26.0	2.88	1.94	2.51	2.67	2.74	1.84	2.37	2.53
	30.0	2.79	1.90	2.44	2.59	2.67	1.80	2.32	2.46
	60.0	2.41	1.74	2.15	2.27	2.27	1.65	2.04	2.14
	90.0	2.20	1.66	2.01	2.09	2.11	1.57	1.91	1.99

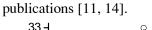
Table 1: The height of first peaks of the site-site distribution functions derived with GEMC-NPT simulation using the *ab initio* pair potentials Eq. 1 and Eq. 2

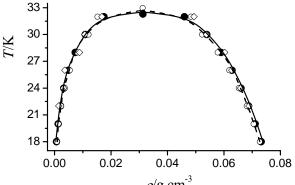
2. Phase coexistence properties

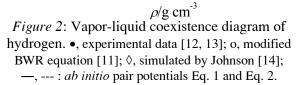
The simulation results are shown in tables 2 and 3. The vapor-liquid coexisting phase curve of the fluid hydrogen is illustrated in figure 2. Experimental data [12, 13], values from the modified empirical equation of state [11] as well as from Johnson simulation results using the Goldman (SG) potential [14] are also included.

The vapor pressures and enthalpies of vaporization derived from the same simulations are depicted in figures 3a and 3b.

These vapor pressures differ on absolute average from the experimental data typically by about 3.49% and 9.14%. These differences are small within statistical uncertainties of experimental resources and a few previous







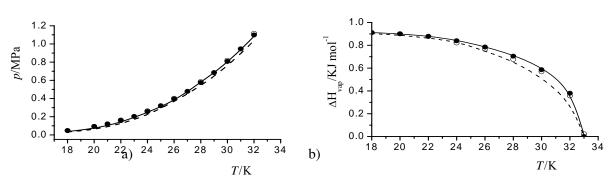


Figure 3: a) Vapor pressure, b) Vaporization enthalpy; for an explanation see Fig. 2

Method	<i>T</i> _c , K	$ ho_{\rm c}, {\rm g.cm}^{-3}$	P _c , MPa	$V_{\rm c},{\rm cm}^3{\rm mol}^{-1}$	Ref.
Eq. 1	33.216	0.0313	1.1258	64.3806	this work
Eq. 2	33.024	0.0311	1.0990	64.7907	this work
EOS	32.972	0.0312	1.2837	64.1539	[11]
Exp.	33.190	0.0312	1.2928	64.1026	[12]
Exp.	33.00	0.0310	1.2930	64.5677	[13]

Table 2: Critical properties of hydrogen resulting from the GEMC-NVT simulation using equations Eq. 1 and 2; EOS: equation of state [11]; Exp.: experimental values

The critical properties of the pure fluid hydrogen could not be calculated directly from the simulation, but they could be obtained from the orthobaric densities of vapor-liquid equilibria by the least-square fit to the relations (3), as shown in table 2. The critical pressure of hydrogen was calculated from the Antoine equation Eq. 4. The results agreed reasonable well with experimental data. The thermodynamic properties of this fluid are also shown in table 3.

Table 3: Enthalpy of vaporization $\Delta_{vap}H$, entropy of vaporization $\Delta_{vap}S$ and boiling temperature T_b at P = 101.3 kPa predicted from simulation vapor pressures

Method	$\Delta_{\rm vap}H$, kJ mol ⁻¹	$\Delta_{\rm vap}S$, kJ.mol ⁻¹ .K ⁻¹	<i>T</i> _b , K	ref.
Eq. 1	1.17148	0.05608	20.8911	this work
Eq. 2	1.21621	0.05717	21.2740	this work
EOS	1.07399	0.05305	20.2457	[11]
Exp.	1.07786	0.05299	20.3900	[12]
Exp.	1.07752	0.05314	20.2754	[13]

The discrepancies between predicted results and experimental data are insignificant.

IV - CONCLUSION

The vapour-liquid phase equilibria and thermodynamic properties of the pure fluid hydrogen were calculated successfully with our developed computer simulation programs GEMC-NVT and GEMC-NPT using *ab initio* intermolecular pair potentials. The simulation results turn out to be in very good agreement with experimental data and with those from literature data. This also confirmed that the two our new 5-site analytical potential functions resulting from *ab initio* calculations in the work [8] are of high quality, accurate and reliable.

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