

## Accelerated Methods for Solving Grid Equations. I

Dang Quang A  
Institute of Informatics  
National Centre for Scientific Research of Vietnam

### 1. Introduction

The boundary value problems for elliptic equations after discretization by the method of finite differences or by the finite elements method are reduced to the large system of linear algebraic equations

$$Au = f. \quad (1.1)$$

where  $A$  is a symmetric, positive definite matrix of order  $N$  equal to the number of nodes of the grid. There have been many effective methods for solving the system (1.1) (see e.g. [1-4]). Besides, in order to accelerate the convergence rate of the main iterative processes one proposed some interesting methods. First we should mention the group of accelerated methods by additive correction ( see [3] and bibliography therein ), whose most important representatives are the multigrid methods [5-7]. These methods are intensively developed now. The second group of accelerated methods are those by multiple correction [8,9]. In spite of the existing methods the elaboration of new methods either more effective or easily realized on parallel computers is of interest to researchers.

In this work we propose a method accelerating the convergence rate of certain well known iterative methods such as the alternately triangular method (ATM) and the alternating direction method (ADM) [1,4,16,17]. More precisely, we suggest new algorithms on the base of the above mentioned methods and prove that our algorithms converge faster than ATM and ADM applied directly to the system (1.1). Our idea is based on the construction of an appropriate combination of the solutions of the systems, which may be solved faster than the original system (1.1). This idea on the continuous level was used by ourselves in [10-12] to construct effective algorithms for solving biharmonic equation and the second order elliptic equation with discontinuous coefficients.

is possible to consider this idea as an expression of the new approach to numerical modelling of problems in physics and mechanics stated recently in [13,14].

### Construction of approximate solution of operator equation by parametric extrapolation

In this section we construct the theoretical background of our accelerated algorithms. First we establish some auxiliary results. Let  $A$  be a linear, symmetric, positive operator in the Hilbert space  $H$  with the scalar product  $(\cdot, \cdot)$ . Consider the equation

$$Au = f, \quad f \in R(A). \quad (2.1)$$

According to this equation we introduce the perturbed equation

$$Au_\epsilon + \epsilon Pu_\epsilon = f \quad (2.2)$$

where  $\epsilon$  is a small parameter,  $\epsilon > 0$ ;  $P$  is a linear, symmetric, positive definite operator  $P = P^* \geq \mu E$ ,  $\mu > 0$ ,  $E$  is the identity operator.

**Lemma 2.1.** Let  $u_0$  and  $u_\epsilon$  be the solutions of the equations (2.1), (2.2) respectively, where  $A$ ,  $P$  are the operators described above. Then

$$(Au_\epsilon, u_\epsilon) \leq (Au_0, u_0) = (f, u_0). \quad (2.3)$$

*Proof.* From (2.1) and (2.2) it follows that

$$Au_\epsilon + \epsilon Pu_\epsilon = Au_0.$$

Multiplying both sides of the above equality by  $u_\epsilon$  and taking into account that  $P > 0$  we get

$$(Au_\epsilon, u_\epsilon) \leq (Au_0, u_\epsilon). \quad (2.4)$$

From the inequality

$$(A(u_\epsilon - u_0), u_\epsilon - u_0) = (Au_\epsilon, u_\epsilon) + (Au_0, u_0) - 2(Au_0, u_\epsilon) \geq 0$$

we derive

$$2(Au_0, u_\epsilon) \leq (Au_\epsilon, u_\epsilon) + (Au_0, u_0)$$

and this estimate and (2.4) imply the desired inequality (2.3).

**Remark 2.1.** In the case when  $H$  is finite-dimensional space from (2.3) we have

$$\|u_\epsilon\|^2 \leq C_1(f, u_0),$$

where  $C_1 = \text{const.}$ , independent of  $\epsilon$ .

**Remark 2.2.** If in addition to the assumptions of Lemma 2.1 we suppose that the operator  $A$  is completely continuous and the operators  $A$  and  $P$  are commute then instead of (2.3) we get the stronger estimate

$$\|u_\epsilon\| \leq \|u_0\|.$$

This result in the case  $P = E$  was proved in [10,12]. Here it should take into account the commutativity of  $A$  and  $P$ .

In the future we shall need the following result.

**Lemma 2.2.** (see [15], Appendix). Let  $M$  be a natural number. Then the solution of the system

$$\sum_{i=1}^{M+1} \gamma_i = 1, \quad \sum_{i=1}^{M+1} \frac{1}{i!} \gamma_i = 0, \quad i = \overline{1, M}$$

is given by the formula

$$\gamma_i = \frac{(-1)^{M+1-i} i^{M+1}}{i!(M+1-i)!}, \quad i = \overline{1, M+1}. \quad (2.5)$$

Now we prove a theorem on asymptotic expansion of the solution of the perturbed equation (2.2).

**Theorem 2.1.** Let  $A$  be a linear, symmetric, positive operator;  $P$  be a linear, symmetric, positive definite operator in the Hilbert space. Let  $M$  be a given natural number such that the equations

$$Av_k = -Pv_{k-1}, \quad k = \overline{1, M+1}, \quad v_0 = u_0 \quad (2.6)$$

have solutions. Then the solution of the equation (2.2) can be expanded in the form

$$u_\epsilon = u_0 + \sum_{k=1}^M \epsilon^k v_k + \epsilon^{M+1} w_\epsilon, \quad (2.7)$$

where  $u_0$  is the solution of (2.1),  $v_k$ , ( $k = \overline{1, M}$ ) are elements independent of  $\epsilon$ , and  $w_\epsilon$  satisfies the estimate

$$(Aw_\epsilon, w_\epsilon) \leq C_2 = \text{const}. \quad (2.8)$$

*Proof.* Substitute (2.7) into (2.2). After balancing the coefficients of  $\epsilon$  we get the equations

$$Av_k = -Pv_{k-1}, \quad k = \overline{1, M},$$

$$(A + \epsilon P)w_\epsilon = -Pv_M.$$

Clearly,  $v_k$  ( $k = \overline{1, M}$ ) do not depend on  $\epsilon$  and by Lemma 2.1 we have

$$(Aw_\epsilon, w_\epsilon) \leq (Aw_0, w_0),$$

where  $w_0$  is the solution of the equation solvable by assumption

$$Aw_0 = -Pv_M.$$

Thus, the theorem is proved with  $C_2 = (Aw_0, w_0)$ .

In order to use Theorem 2.1 for constructing solution of grid equations we restate one in the case of finite-dimensional space with the attention that in this space positive operator is also positive definite.

**Theorem 2.1a.** Let  $A$  and  $P$  be linear, symmetric operators in  $N$ -dimensional Euclidean space. Moreover,  $A > 0$ ,  $P \geq 0$ . Then for any natural number  $M$  it is possible to expand the solution of the equation (2.2) in the form

$$u_\epsilon = u_0 + \sum_{k=1}^M \epsilon^k v_k + \epsilon^{M+1} w_\epsilon, \quad (2.9)$$

where  $v_k$  ( $k = \overline{1, M}$ ) are elements independent of  $\epsilon$ ,

$$\|w_\epsilon\| \leq C_3. \quad (2.10)$$

We now construct the solution of (2.1) by extrapolation of solutions of (2.2).

Put

$$U^E = \sum_{k=1}^{M+1} \gamma_k u_{\epsilon/k}, \quad (2.11)$$

where  $\gamma_k$  ( $k = \overline{1, M+1}$ ) are coefficients given by (2.5),  $u_{\epsilon/k}$  are the solutions of (2.2) with the parameters  $\epsilon/k$ . These solutions are called the basic solutions and  $U^E$  is considered as an approximate solution of the equation (2.1).

**Theorem 2.2.** Let  $A$  and  $P$  be operators, satisfying the assumptions of Theorem 2.1a. Then for the approximate solution of (2.1) constructed by (2.11) there holds the estimate

$$\|U^E - u_0\| \leq C_4 \epsilon^{M+1}. \quad (2.12)$$

*Proof.* We use the expansion (2.9) for  $u_{\epsilon/k}$  ( $k = \overline{1, M+1}$ ) and form  $U^E$  by (2.11). Then applying Lemma 2.1 and the estimate (2.10) we obtain (2.12) with

$$C_4 = C_3 \sum_{k=1}^{M+1} \frac{|\gamma_k|}{k^{M+1}}.$$

Thus, the theorem is proved.

We see from above that in order to construct approximate solution of the original problem (2.1) we have to solve a number of perturbed problems (2.2) with the parameters  $\epsilon/k$  ( $k = \overline{1, M+1}$ ). Usually, we can not find exact solutions of these problems but only their approximate solutions. Suppose that by iterative method we have found approximate solutions  $y_{\epsilon/k}^{(n_k)}$  of the perturbed equations with accuracy  $\theta$ , i.e.

$$\|y_{\epsilon/k}^{(n_k)} - u_{\epsilon/k}\| \leq \theta, \quad k = \overline{1, M+1}. \quad (2.13)$$

Then instead of the theoretical approximate solution  $U^E$  we obtain the real one

$$Y^E = \sum_{k=1}^{M+1} \gamma_k y_{\epsilon/k}^{(n_k)}. \quad (2.14)$$

It is easy to get



Lemma 2.3. For the real approximate solution  $Y^E$  we have

$$\|Y^E - u_0\| \leq C_4 \epsilon^{M+1} + C_5 \theta,$$

where

$$C_5 = \sum_{k=1}^{M+1} |\gamma_k|.$$

Let  $M$  be chosen, that is we have fixed the number of basic solutions for constructing  $U^E$ . From Lemma 2.3 we obtain

Theorem 2.3. If  $\epsilon^{M+1} = \theta$  then  $Y^E$  computed by (2.14) is approximate solution of the problem (2.1) with the accuracy

$$\|Y^E - u_0\| \leq C\theta,$$

where  $C = C_4 + C_5$  independent of  $\theta$ .

Suppose that the equations (2.1), (2.2) are solved by such a same iterative method. Then the construction of the solution of (2.1) by the parametric extrapolation (2.11) is effective if the total computational cost needed for solving the perturbed equations is less than for solving the original equation. In Section 3 this fact will be established for some iterative methods.

Remark 2.3. The basic solutions  $u_{\epsilon/k}$  are sought independently from each other. Hence, it is optimal to realize the parametric extrapolation method on parallel computers. Otherwise, on sequential computers if the perturbed equations are solved iteratively then for reducing the computation time it is proposed to take  $\frac{1}{k} \sum_{i=1}^k y_{\epsilon/i}^{(n_i)}$  to be the starting approximation in the iterative process for finding  $u_{\epsilon/(k+1)}$ . Here  $y_{\epsilon/i}^{(n_i)}$  is the approximation of  $u_{\epsilon/i}$  with given accuracy.

### 3. Application to grid equations

In the  $N$ -dimensional Euclidean space consider the equation

$$Au = f, \quad (3.1)$$

where  $A$  is a linear, symmetric, positive definite operator:

$$A = A^* \geq \delta E, \quad \delta > 0. \quad (3.2)$$

For solving the above equation one usually constructs two-layer iterative process of the form

$$B \frac{y^{(k+1)} - y^{(k)}}{\tau_{k+1}} + Ay^{(k)} = f, \quad k = 0, 1, \dots, n-1. \quad (3.3)$$

Here  $B$  is an operator, satisfying the following requirements: i)  $B$  may be fast inverted ii) the ratio of the energetic equivalence coefficients of the operators  $A$  and  $B$  reaches maximal value, i.e.  $\gamma_1/\gamma_2 \rightarrow \max$ , where  $\gamma_1, \gamma_2$  are the coefficients in the relation

$$\gamma_1 B \leq A \leq \gamma_2 B,$$

parameters of the Chebyshev collection (see [1,16]). The convergence rate depends on  $1/\gamma_2$ .

Usually, for the construction of  $B$  one starts from an operator  $R = R^* > 0$ , which is energetic equivalent to  $A$ , i.e.

$$C_1 R \leq A \leq C_2 R, \quad C_2 \geq C_1 > 0, \quad (3.4)$$

Using the above idea and the results of Section 2 we shall construct efficient algorithm for solving equation (3.1).

Suppose that the operator  $R$  can be decomposed into the sum of operators. We shall consider following cases.

**Case 1:**  $R$  is the sum of two noncommute but conjugate each other operators i.e.

$$R = R_1 + R_2, \quad R_2^* = R_1, \quad R_1 R_2 \neq R_2 R_1.$$

In this case we choose  $P = R_1 R_2$ . Clearly,  $P = P^* > 0$ .

Consider the perturbed equation

$$A_\epsilon u_\epsilon = f, \quad (3.5)$$

where

$$A_\epsilon = A + \epsilon P, \quad \epsilon > 0. \quad (3.6)$$

We shall construct an operator  $B$ , energetic equivalent to  $A_\epsilon$  in the form

$$B = (E + \omega R_1)(E + \omega R_2). \quad (3.7)$$

The problem is to choose  $\omega$  so that the ratio  $\gamma_1/\gamma_2$  reaches the more the better value. Here  $\gamma_1, \gamma_2$  are the coefficients in the relation

$$\gamma_1 B \leq A_\epsilon \leq \gamma_2 B \quad (3.8)$$

**Lemma 3.1.** Let us choose  $\omega = a\sqrt{\epsilon}$ , where  $a$  is a real number, satisfying

$$a \geq \frac{\sqrt{\epsilon}}{C_1}, \quad (3.9)$$

being the coefficient in (3.4).

Then we have the inequalities (3.8), where

$$\gamma_1 = \frac{\delta}{1 + \delta a^2}, \quad \gamma_2 = \frac{C_2}{a\sqrt{\epsilon}}. \quad (3.10)$$

If  $\epsilon < C_1^2/\delta$  then the ratio  $\xi = \gamma_1/\gamma_2$  reaches maximum when

$$a = a^0 = 1/\sqrt{\delta} \quad (3.11)$$

In this case we have

$$\xi = \overset{\circ}{\xi} = \frac{\sqrt{\delta}}{2C_2} \sqrt{\epsilon}, \quad (3.1)$$

$$\gamma_1 = \overset{\circ}{\gamma}_1 = \delta/2, \quad \gamma_2 = \overset{\circ}{\gamma}_2 = C_2 \sqrt{\frac{\delta}{\epsilon}}. \quad (3.1)$$

*Proof.* For all  $x \in R^N$  we have

$$(Bx, x) = \|x\|^2 + \omega(Rx, x) + \omega^2 \|R_2x\|^2,$$

$$(A_\epsilon x, x) = (Ax, x) + \epsilon \|R_2x\|^2.$$

Inserting  $\omega = a\sqrt{\epsilon}$  into the above equalities and taking into account (3.4) we get

$$\begin{aligned} (Bx, x) &\geq \|x\|^2 + \frac{a\sqrt{\epsilon}}{2}(Ax, x) + a^2\epsilon \|R_2x\|^2 \\ &\geq \frac{a\sqrt{\epsilon}}{C_2} ((Ax, x) + C_2 a\sqrt{\epsilon} \|R_2x\|^2). \end{aligned}$$

Using (3.9) we obtain

$$(Bx, x) \geq \frac{a\sqrt{\epsilon}}{C_2} ((Ax, x) + \epsilon \|R_2x\|^2) = \frac{a\sqrt{\epsilon}}{C_2} (A_\epsilon x, x).$$

that is

$$B \geq \frac{a\sqrt{\epsilon}}{C_2} A_\epsilon. \quad (3.1)$$

Next we get an upper bound for  $B$ . We have

$$\begin{aligned} (Bx, x) &\leq \|x\|^2 + \frac{a\sqrt{\epsilon}}{C_1}(Ax, x) + a^2\epsilon \|R_2x\|^2 \\ &\leq \frac{1}{\delta}(Ax, x) + a^2 \left( \frac{\sqrt{\epsilon}}{aC_1}(Ax, x) + \epsilon \|R_2x\|^2 \right) \\ &\leq \frac{1}{\delta}(A_\epsilon x, x) + a^2((Ax, x) + \epsilon \|R_2x\|^2) \leq \left( \frac{1}{\delta} + a^2 \right) (A_\epsilon x, x). \end{aligned}$$

Thus, we get

$$B \leq \left( \frac{1}{\delta} + a^2 \right) A_\epsilon. \quad (3.1)$$

From (3.14), (3.15) we obtain (3.8) with  $\gamma_1, \gamma_2$  computed by (3.10). Their ratio is a function of

$$\xi = \xi(a) = \frac{\gamma_1}{\gamma_2} = \frac{\delta\sqrt{\epsilon}a}{C_2(1 + \delta a^2)}.$$

If  $0 < \epsilon < C_1^2/\delta$  then  $\max_{a \geq \sqrt{\epsilon}/C_1} \xi(a)$  is reached when  $a = \overset{\circ}{a} = 1/\sqrt{\delta}$ . In this case we have  $\xi = \overset{\circ}{\xi}, \gamma_1 = \overset{\circ}{\gamma}_1, \gamma_2 = \overset{\circ}{\gamma}_2$  computed by (3.12), (3.13).

The lemma is proved.

For solving the perturbed equation we use the iterative process

$$B \frac{y_\epsilon^{(k+1)} - y_\epsilon^{(k)}}{\tau_{k+1}} + A_\epsilon y_\epsilon^{(k)} = f, \quad k = \overline{0, n-1}, \quad (3.16)$$

$y_\epsilon^{(0)}$  is given.

From the theory of two-layer iterative processes in [1,16] we obtain the following

**Theorem 3.1.** *If*

$$B = (E + \sqrt{\frac{\epsilon}{\delta}} R_1)(E + \sqrt{\frac{\epsilon}{\delta}} R_2)$$

and  $\{\tau_k\}_{k=1}^n$  is the Chebyshev parameters collection, constructed according to the bounds  $\gamma_1, \gamma_2$  defined in Lemma 3.1 then for the iterative process (3.16) there holds the estimate

$$\|y_\epsilon^{(n)} - u_\epsilon\| \leq q_n \|y_\epsilon^{(0)} - u_\epsilon\|, \quad (3.17)$$

where

$$q_n = \frac{2\rho_1^n}{1 + \rho_1^{2n}}, \quad \rho_1 = \frac{1 - \sqrt{\xi}}{1 + \sqrt{\xi}}, \quad \xi = \frac{\sqrt{\delta}\sqrt{\epsilon}}{2C_2} \quad (3.18)$$

in order to achieve the relative accuracy  $\theta$ , that is to have

$$\|y_\epsilon^{(n)} - u_\epsilon\| \leq \theta \|y_\epsilon^{(0)} - u_\epsilon\| \quad (3.19)$$

is needed to implement  $n$  iterations,  $n \geq n_c(\theta)$ ,

$$n_c(\theta) = \ln \frac{2}{\theta} / \ln \frac{1}{\rho_1} \sim \ln \frac{2}{\theta} / (\sqrt{\frac{2}{C_2}} \delta^{1/4} \epsilon^{1/4}).$$

For the stationary iterative process with

$$\tau_k \equiv \tau = \frac{2}{\gamma_1 + \gamma_2}$$

instead of (3.17) we have

$$\|y_\epsilon^{(k)} - u_\epsilon\| \leq \rho^k \|y_\epsilon^{(0)} - u_\epsilon\|, \quad k = 1, 2, \dots$$

where

$$\rho = \frac{1 - \xi}{1 + \xi}$$

In this case the number of iterations needed to achieve the relative accuracy  $\theta$  is  $n$ ,  $n \geq n_d(\theta)$ ,

$$n_d(\theta) = \ln \frac{2}{\theta} / \left( \frac{\sqrt{\delta}}{C_2} \sqrt{\epsilon} \right).$$

**Example 3.1** We illustrate the proposed method on the Dirichlet problem for the second order elliptic equation

$$\sum_{\alpha=1}^p \frac{\partial}{\partial x_{\alpha}} \left( k_{\alpha}(x) \frac{\partial u}{\partial x_{\alpha}} \right) = -f(x), \quad x \in \Omega, \quad (3.22)$$

$$u = 0, \quad x \in \Gamma.$$

Here  $\bar{\Omega}$  is a parallelepiped in  $p$ -dimensional Euclidean space ( $p = 2, 3$ )

$$\bar{\Omega} = \{0 \leq x_{\alpha} \leq l_{\alpha}, \alpha = \overline{1, p}\},$$

$\Gamma$  is the boundary of  $\bar{\Omega}$ .

We introduce over  $\bar{\Omega}$  the grid

$$\bar{\omega} = \{x_{\alpha} = i_{\alpha} h_{\alpha}, 0 \leq i_{\alpha} \leq N_{\alpha}, \alpha = \overline{1, p}\}$$

with the boundary  $\gamma$ . Here  $N_{\alpha} h_{\alpha} = l_{\alpha}, \alpha = \overline{1, p}$ .

We approximate the problem (3.22) by the difference scheme

$$\Delta y \equiv \sum_{\alpha=1}^p (a_{\alpha} y_{\bar{x}_{\alpha}})_{x_{\alpha}} = -f(x), \quad x \in \omega, \quad (3.23)$$

$$y = 0 \quad x \in \gamma.$$

Here and henceforward we use the difference notations taken from [1,16].

Let us denote by  $\overset{0}{H}_h$  the space of grid functions defined on  $\bar{\omega}$  and vanishing on  $\gamma$ . Then the difference scheme may be written in the form of operator equation

$$Ay = f. \quad (3.24)$$

where  $Ay = -\Delta y, y \in \overset{0}{H}_h$ .

In  $\overset{0}{H}_h$  we also define the operators  $R, R_1, R_2$  as follows

$$R_1 y = \sum_{\alpha=1}^p \frac{y_{\bar{x}_{\alpha}}}{h_{\alpha}}, \quad R_2 y = - \sum_{\alpha=1}^p \frac{y_{x_{\alpha}}}{h_{\alpha}}, \quad R y = \sum_{\alpha=1}^p y_{\bar{x}_{\alpha} x_{\alpha}}.$$

If in (3.22) it is assumed that

$$C_1 \leq k_{\alpha}(x) \leq C_2, \quad C_2 \geq C_1 > 0, \quad \alpha = \overline{1, p}$$

then we have

$$C_1 R \leq A \leq C_2 R.$$

For  $\delta$  in (3.2) we can take

$$\delta = 8 \sum_{\alpha=1}^p \frac{1}{l_{\alpha}^2}.$$

Since the difference scheme (3.23) of the problem (3.22) has accuracy of order  $h^2$ , ( $h^2 = \sum_{\alpha=1}^p h_{\alpha}^2$ ), in order to have consistence we also implement iterative processes with accuracy  $\theta = h^2$ . With this accuracy when fixing  $M = 1$ , i.e. when solving two perturbed problems it should be taken  $\epsilon = h$ . Then the perturbed problem (2.2) has the form

$$Ay + hR_1R_2y = f.$$

It is easy to verify that the number of iterations needed for solving this problem with accuracy  $\theta = h^2$  is

$$n_c(\theta) = \mathcal{O}\left(\frac{1}{h^{1/4}} \ln \frac{1}{h}\right).$$

Hence, the total number of iterations for obtaining the approximate solution of the problem (3.22) with the same accuracy is a quantity of order  $\mathcal{O}\left(\frac{1}{h^{1/4}} \ln \frac{1}{h}\right)$ . This number is less than  $\mathcal{O}\left(\frac{1}{h^{1/2}} \ln \frac{1}{h}\right)$ , which the alternately triangular method with the Chebyshev parameters collection requires (see [1,16]) if applying directly this method to the original difference problem (3.24).

When  $M > 1$  we must solve more than two perturbed problems. In this case, choosing  $\epsilon = h^{2/(M+1)}$  we shall reduce the number of iterations needed for obtaining the approximate solution of (3.22) with accuracy  $h^2$  to  $\mathcal{O}\left(\frac{1}{h^{1/2(M+1)}} \ln \frac{1}{h}\right)$ .

**3.2 Case 2 :**  $R$  is decomposed into the sum of two symmetric, commute operators

$$R = R_1 + R_2, \quad R_{\alpha}^* = R_{\alpha} \quad (\alpha = 1, 2), \quad R_1R_2 = R_2R_1.$$

Suppose  $\delta_{\alpha}E \leq R_{\alpha} \leq \Delta_{\alpha}E$ ,  $\delta_{\alpha} > 0$ ,  $\alpha = 1, 2$ .

Taking  $P = R_1R_2$  and making the same arguments as in Case 1 we obtain similar results.

**Example 3.2.** Consider the Dirichlet problem (3.22) with  $p = 2$ . We also approximate it by the difference scheme (3.23). The operators  $R$ ,  $R_1$ ,  $R_2$  are defined as follows

$$Ry = -y_{\bar{x}_1x_1} - y_{\bar{x}_2x_2}, \quad R_1y = -y_{\bar{x}_1x_1}, \quad R_2y = -y_{\bar{x}_2x_2}$$

The estimate for the number of iterations is similar to that in Example 3.1, where  $R$  is the sum of two operators conjugate each other. Comparing with ADM applied directly to the original difference problem we see the effectiveness of our parametric extrapolation technique.

The case, where  $R$  is split into the sum of three commute operators, which we meet when solving elliptic problems in three-dimensional domains and a comparative analysis of the proposed algorithms on numerical examples will be presented in an another paper.

This work is completed with the financial support from the National Basic Research Program in Natural Sciences.

## References

1. Samarskii A.A. & Nikolaev N.C. Numerical methods for grid equations. M.: Nauka, 1978 (in Russian); Birkhauser Verlag, 1989.
2. Marchuk G.I. Methods of computational mathematics. M.: Nauka, 1989 (in Russian).
3. Hageman L.A., Young D.M. Applied Iterative Methods. Academic Press, 1981.
4. Samarskij A.A. and ets. Certain modern methods for solving grid equations. *Izv. VUZov, Matematika*, 1983, No 7, pp. 3-13 (in Russian).
5. Hackbush W. Multigrid methods and Applications. Springer Ser. Comput. Math. 4, Springer Berlin, 1985.
6. Proceedings of the Fourth Copper Mountain Conference on Multigrid Methods. Edited by Jan Mandel, Stephen F. Mc. Cormic, 1989.
7. Shaidurov V.V. Multigrid finite elements methods. M.: Nauka, 1989 (in Russian).
8. Nakamura S. Effect of weighting functions on the coarse mesh rebalancing acceleration. Proc. Conf. Math. Models and Comput. Tech. for Anal. Nucl. Syst. Vol. II, pp. 120-135. US Atomic Energy Commission, 1974.
9. Wachspress E.L. Iterative Solution of Elliptic Systems and Applications to the Neutron Diffusion Equation. Prentice Hall, Englewood Cliffs, New Jersey, 1966.
10. Dang Quang A. Application of extrapolation to constructing effective method for solving the Dirichlet problem for biharmonic equation. Institute of Computer Science. Preprint 5/1990 (submitted to Differential equations).
11. Dang Quang A. Numerical method for solving the Dirichlet problem for a fourth order equation. Proceedings of the third national conference on gas and fluid mechanics (Hanoi, 1990). Hanoi, 1991, pp. 195-199 (in Vietnamese).
12. Dang Quang A. Approximate method for solving an elliptic problem with discontinuous coefficients. *Journal of Computational and Applied Mathematics*, 1993 (to appear).
13. Belotserkovskii O.M., Tshennikov V.V. Principles of rational numerical modelling in aerohydrodynamics. In book: Rational numerical modelling in nonlinear mechanics. M.: Nauka, 1990 (in Russian).
14. Belotserkovskii O.M., Skobelev B.Ju., Tshennikov V.V. New principles of difference discretization of mathematical physics problems. In the same book.
15. Marchuk G.I., Chaidurov V.V. Raising accuracy of the solutions of difference schemes. M.: Nauka, 1979 (in Russian).
16. Samarskii A.A. Introduction into the theory of difference schemes. M.: Nauka, 1971 (in Russian).

**Abstract.** *In this paper we propose a technique accelerating the convergence rate of the known iterative schemes for solving grid equations such as the alternately triangular method and the alternating direction method. Our idea is by the parametric extrapolation of the solutions of equations, which can be solved faster than the original ones. The efficiency of the accelerated methods is shown on examples.*