

OPTO-ELECTRONIC PROPERTIES OF SMALL ZnO₂ NANOPARTICLES: FIRST-PRINCIPLES INSIGHTS

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Abstract. *We performed density functional theory based first-principles calculations to investigate the opto-electronic properties of small ZnO₂ nanoparticles. We have shown that these ZnO₂ nanoparticles can exhibit either semiconductor or metallic properties depending on the morphology and sizes of the nanoparticles defined by Miller indices. The absorption spectra computed from the real and imaginary parts of the dielectric functions demonstrate that these nanoparticles exhibit rich optical features with highly spatial anisotropy in the visible light range, suggesting that these newly obtained configurations strongly influence the electronic structures of ZnO₂ nanoparticles. Our results propose the potential application of using Zinc peroxide nanoparticles as prospective building blocks for developing early diagnosis nanodevices in drug industry.*

Keywords: density functional theory, opto-electronic properties, Zinc peroxide, nanoparticles.

Classification numbers: 31.15.E; 42.82.Fv.

I. INTRODUCTION

Nano-scale materials, a rich class of nano-particles, nano-wires, nanoribbons, thin films, and component structures, have received a vast attention from both theoretical and experimental physicists since the discovery of fullerenes and the emergence of carbon nanotubes in the 1990s [1–6]. Among these types of nano-scale materials, nano-particles have been strongly investigated in multidisciplinary industries, including medicine, cosmetics, materials, energy and environment [7–13]. Nano-particles with a size of smaller than 100 nm existing in multiple structures varying from zero to three-dimensions [14, 15] are the core strategy to improve the bioavailability, substance stability and safety of the next generation of drug technology [16, 17]. It is because these nanoparticles, thanks to their small sizes, large surface free energies and large contact areas, can significantly magnify the dissolution rate, thereby increase the bioavailability of drugs in living organisms. Furthermore, under the same activation conditions, some targeted nanoparticles sufficiently reinforce the absorption efficiency of drugs via their ability to shine 20 times brighter and maintain the brightness coefficient much longer than fluorescent molecules [16]. The ultimate combination of highly mobile and long-lived photoluminescence of targeted nano-particles allows physicians to follow living cells deep inside the body with better image contrast, and so as undoubtedly useful for early diagnosis [18].

Among the class of nano-scale materials, recently, TiO₂, ZnO, and ZnO₂ have been shown as excellent candidates for their potential applications in photocatalytics [19], gas sensing [20], photovoltaics [21] and especially in pharmaceutical technology thanks to their antibacterial efficacies [22, 23].

More interestingly, ZnO₂ (Zinc peroxide), a novel material with rich morphological configurations, promises great applications in (pharmaceutical-) industry. ZnO₂ is thermally stable up to 250 °C and can be created in vast of synthesized methods [24–28]. Also, ZnO₂ can be hydrothermalized from a sol-gel phase to be used as a precursor to create nanometric ZnO particles [22, 29] and ultrasmall functional ZnO₂ nanoparticles [30].

While detailed studies on the electronic and optical properties of ZnO nanoparticles have been reported tremendously over the last decade, experimental and theoretical research on the opto-electronic properties of ZnO₂ nanoparticles are still limited. Thus, in this paper, we report the electronic and optical properties of small ZnO₂ nanoparticles with different nanoparticle sizes and shapes (crystal growing planes) using density functional theory. We designed six types of nanoparticles including ZnO₂(001)-21, ZnO₂(001)-41, ZnO₂(001)-187, ZnO₂(111)-13, ZnO₂(111)-33 and ZnO₂(111)-173, where (001) and (111) denote the Miller indices while the latter numbers correspond to the number of Zn and O atoms in these nanoparticles, and examined their electronic and optical properties in detail.

II. COMPUTATIONAL METHODS

In this work, we created ZnO₂ nanoparticles according to the Wulff construction with suitable Miller indices and their corresponding minimum surface energies using the novel web-base crystallographic Nano-Crystal tool [31]. Using this tool, we have found the two most stable facet ZnO₂ nanoparticles with their minimum surface energy values of 62 eV for ZnO₂(001) and 35 eV for ZnO₂(111), respectively. We then created six ZnO₂ samples with increasing nanoparticle sizes from a ZnO₂ pyrite crystal as shown in Fig. 1. The ZnO₂ crystallographic file can be downloaded

from Ref. [32]. Six investigated samples include three samples in the (001) Miller plane (namely ZnO₂(001)-21, ZnO₂(001)-41, ZnO₂(001)-187, the latter numbers in each sample names denote the total numbers of atoms in each nanoparticles corresponding to the values of 21, 41 and 187, respectively) and three samples in the (111) Miller plane (namely ZnO₂(111)-13, ZnO₂(111)-33 and ZnO₂(111)-173 with the equivalent numbers of atoms of 13, 33 and 173, respectively).

All opto-electronic properties of ZnO₂ nanoparticles were calculated using the density functional theory (DFT) predictions implemented in the Atomistix Toolkit (ATK) package [33, 34]. For the Generalized Gradient Approximation (GGA), the Perdew–Burke–Ernzerhof (PBE) with a double- ζ polarized basis is applied to handle the exchange-correlation function. Monkhorst–Pack grids of $5 \times 5 \times 5$ k-points were used for the geometry optimization, and one k-point (Γ point) for the calculation of optoelectronic properties. The cut-off energy was about 500 eV. Vacuum spaces of at least 15 Å are applied on all three axes of the investigated samples to avoid the interactions between the periodic images. This approach is to ensure the reproducibility and reliability of our opto-electronic calculations for all ZnO₂ nanoparticles.

The optical properties of materials were determined by a complex dielectric function [35, 36]:

$$\varepsilon(\omega) = \varepsilon_1(\omega) + i\varepsilon_2(\omega). \quad (1)$$

Here, $\varepsilon_1(\omega)$ and $\varepsilon_2(\omega)$ are the real part and the imaginary part of the dielectric function, respectively. The dielectric function is related to the susceptibility as

$$\varepsilon(\omega) = 1 + \chi(\omega). \quad (2)$$

where ω indicates the angular frequency of the incident photon. $\chi(\omega)$ is calculated by the Kubo-Greenwood formula [37]. From here, we can determine the extinction coefficient κ , and the optical absorption coefficient α of the materials as follows [35, 38]:

$$\kappa = \sqrt{\frac{\sqrt{\varepsilon_1^2 + \varepsilon_2^2} - \varepsilon_1}{2}}, \quad (3)$$

$$\alpha = \frac{2 \cdot \omega}{c} \kappa, \quad (4)$$

where c is the velocity of light in vacuum. Note that we use the independent particle approximation for linear optical response functions to establish a link between the electronic band structures and the optical properties of the nanoparticles.

III. RESULTS AND DISCUSSION

III.1. Structural properties

Figure 1 and Table 1 present the most energetically favorable ZnO₂ nanoparticle configurations and their size parameters in the edges and the diagonal directions. The distances d_1 , d_2 represent for the sizes of the cubic nanoparticles in the edge and the diagonal directions before DFT optimization while d'_1 , d'_2 are those equivalent lengths after optimization. It is obvious that after optimization, the nanoparticles rearrange in such a way to increase the compactness of the particles resulting in the decrease of d'_1 , d'_2 with respect to d_1 , d_2 distances.

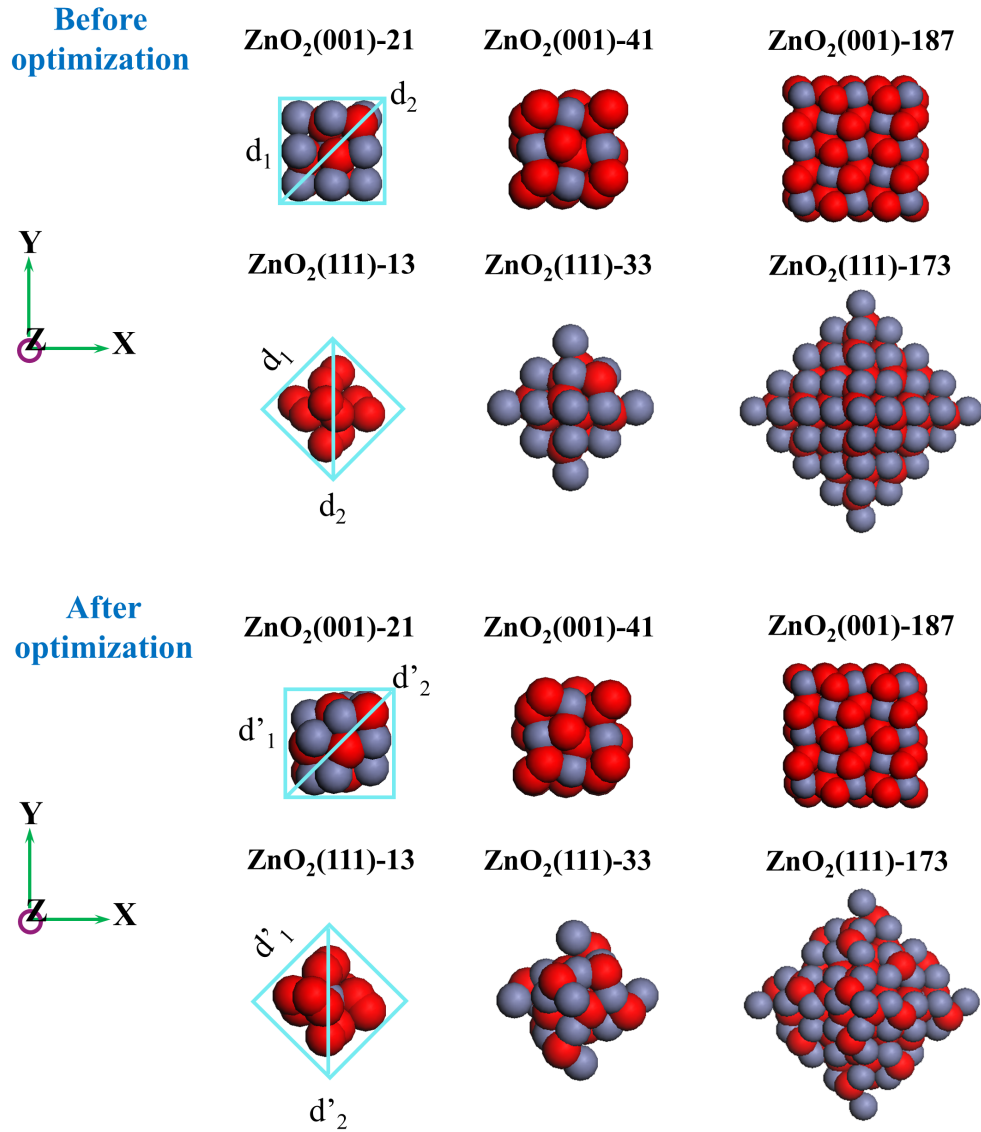


Fig. 1. Configurations of ZnO₂ nanoparticles before and after optimization with different numbers of atoms. Gray and red balls represent for Zinc and Oxygen atoms, respectively.

To confirm that these acquired nanoparticles are the most stable configurations, we calculate their binding energies as following:

$$E_B = \frac{E_{\text{ZnO}_2} - n_{\text{Zn}}E_{\text{Zn}} - n_{\text{O}}E_{\text{O}}}{n_{\text{Zn}} + n_{\text{O}}} \quad (5)$$

Here, E_{ZnO_2} is the total energy of ZnO₂ nanoparticles, E_{Zn} , E_{O} are the energies of the isolated Zinc and Oxygen atoms, and n_{Zn} , n_{O} are the atom numbers of Zinc and Oxygen atoms, respectively.

Table 1. Size parameters in the edges and diagonal directions of the ZnO₂ nanoparticles before and after geometrical optimization.

Samples	d_1 (Å)	d'_1 (Å)	d_2 (Å)	d'_2 (Å)
ZnO ₂ (001)-21	4.90	4.20	6.93	5.90
ZnO ₂ (001)-41	4.96	4.80	7.01	6.74
ZnO ₂ (001)-187	9.92	9.41	14.02	13.30
ZnO ₂ (111)-13	4.21	3.96	5.95	5.68
ZnO ₂ (111)-33	7.01	6.63	9.92	9.45
ZnO ₂ (111)-173	14.02	14.01	19.83	19.83

Table 2. Binding energy of ZnO₂ nanoparticles.

Samples	Zn atoms	O atoms	Binding energy, E_B (eV)
ZnO ₂ (001)-21	13	8	-2.6
ZnO ₂ (001)-41	13	28	-3.7
ZnO ₂ (001)-187	63	124	-3.9
ZnO ₂ (111)-13	1	12	-4.1
ZnO ₂ (111)-33	19	14	-3.1
ZnO ₂ (111)-173	85	88	-2.9

Table 2 shows the binding energies of two groups of ZnO₂(001) and ZnO₂(111) nanoparticles. The calculated binding energy varying in between -2.6 eV and -4.1 eV demonstrates that these investigated nanoparticles are stabilized. For the ZnO₂(001) nanoparticles, their binding energies decrease with increasing particle sizes, suggesting that the larger ZnO₂ nanoparticles are more stable than the smaller ones. However, for ZnO₂(111) nanoparticles, this quantity increases when increasing the nanoparticle sizes, pointing out that larger ZnO₂(111) nanoparticles may be less stable.

III.2. Electronic properties

We now turn our attention to the electronic properties of the optimized ZnO₂ nanoparticles. Figs. 2 and 3 present the calculated band structures and the corresponding density of states of all investigated nanoparticles. For the sake of comparison, we set the Fermi level to zero and only plotted the calculated energy levels 2 eV around the Fermi level. In Fig. 2, it is obvious that ZnO₂ nanoparticles exhibit semiconductor behavior with a bandgap of ~ 0.3 eV for both ZnO₂(001)-21 and ZnO₂(111)-13, while the rest of the investigated samples are electrically metallic. Here, we defined the energy band gap as the energy difference between the highest occupied molecular

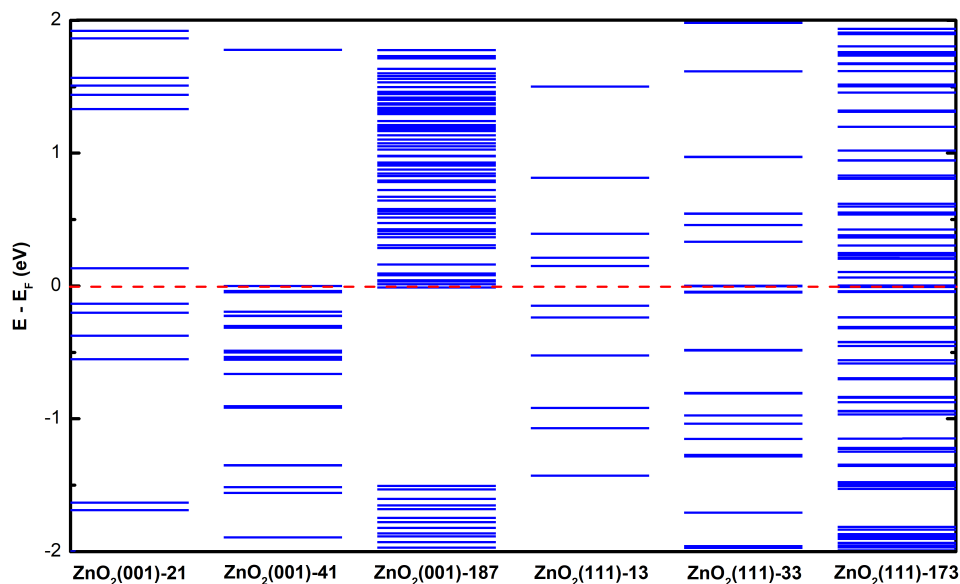


Fig. 2. Energy levels of ZnO_2 nanoparticles, for different sizes and configurations. The dashed lines indicate the Fermi level.

orbitals (HOMO) and the lowest unoccupied molecular orbitals (LUMO). This band structure explicitly indicates that the energy band gaps of the nanoparticles significantly shrink and even vanish when these nanoparticles are big in their atom sizes, pointing out a distinct difference between these nanoparticles and the ZnO_2 pyrite crystal. Note that the band gap of ZnO_2 pyrite crystal is 2.2 eV [32]. Thus, the morphology of the ZnO nanoparticles is central to control the electronic properties of these interesting nanoparticles.

To understand this striking observation, we further explore the role of each type of atoms on the partial density of states in Fig. 3. In the case of ZnO_2 pyrite crystal, Oxygen atoms are the major contributor to both conduction and valence states [32]. However, for $\text{ZnO}_2(001)$ -41, $\text{ZnO}_2(001)$ -187 and $\text{ZnO}_2(111)$ -13, the valence band states are mainly contributed by Oxygen atoms while for the other nanoparticles, Zinc atoms seem to dominate. In the conduction bands, for $\text{ZnO}_2(001)$ -21, $\text{ZnO}_2(001)$ -41, $\text{ZnO}_2(111)$ -33 and $\text{ZnO}_2(111)$ -173, Zinc atoms play a predominant role. An opposite trend was observed for $\text{ZnO}_2(001)$ -187 and $\text{ZnO}_2(111)$ -13. In short, the contribution of Zinc and Oxygen atoms to the density of states varies depending on the surface morphology of certain nanoparticles.

III.3. Optical properties

In this section, we present the optical properties of ZnO_2 nanoparticles. These properties, which are closely related to their electronic band structures, indicate the possibility for using this class of materials for early diagnosis nano-devices in drug industry. To examine the absorption capacity of these nanoparticles, we first calculate the real and imaginary parts of the optical dielectric response functions, and then determine the absorption coefficient ($\alpha(\omega)$) using the equations presented in Sec. II. In these calculations, we have chosen the polarization direction parallel to the electric field.

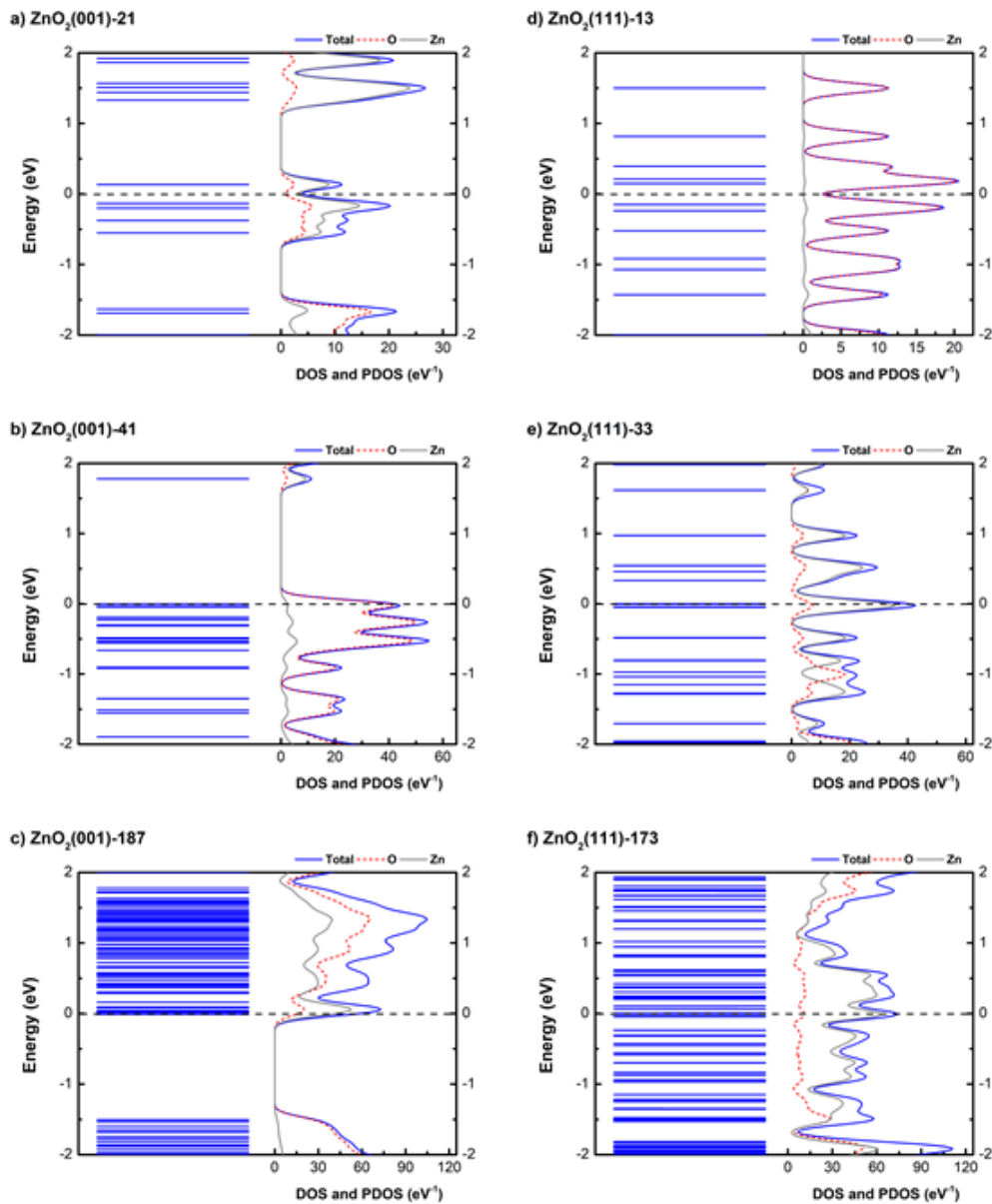


Fig. 3. DOS and PDOS of the ZnO₂ nanoparticles. The dashed lines indicate the Fermi level. Electronic band structure is also included for ease of analysis

The electronic polarizability indicating the absorption capacity of the material can be determined from the real and imaginary parts of the dielectric functions. In Fig. 4, we show the real parts of the dielectric functions of all investigated samples as a function of the incident photon energy. It is clear that for nanoparticles with the number of atoms above 33, the real parts of the dielectric functions do not show a significant difference in all three polarizations. However, for the

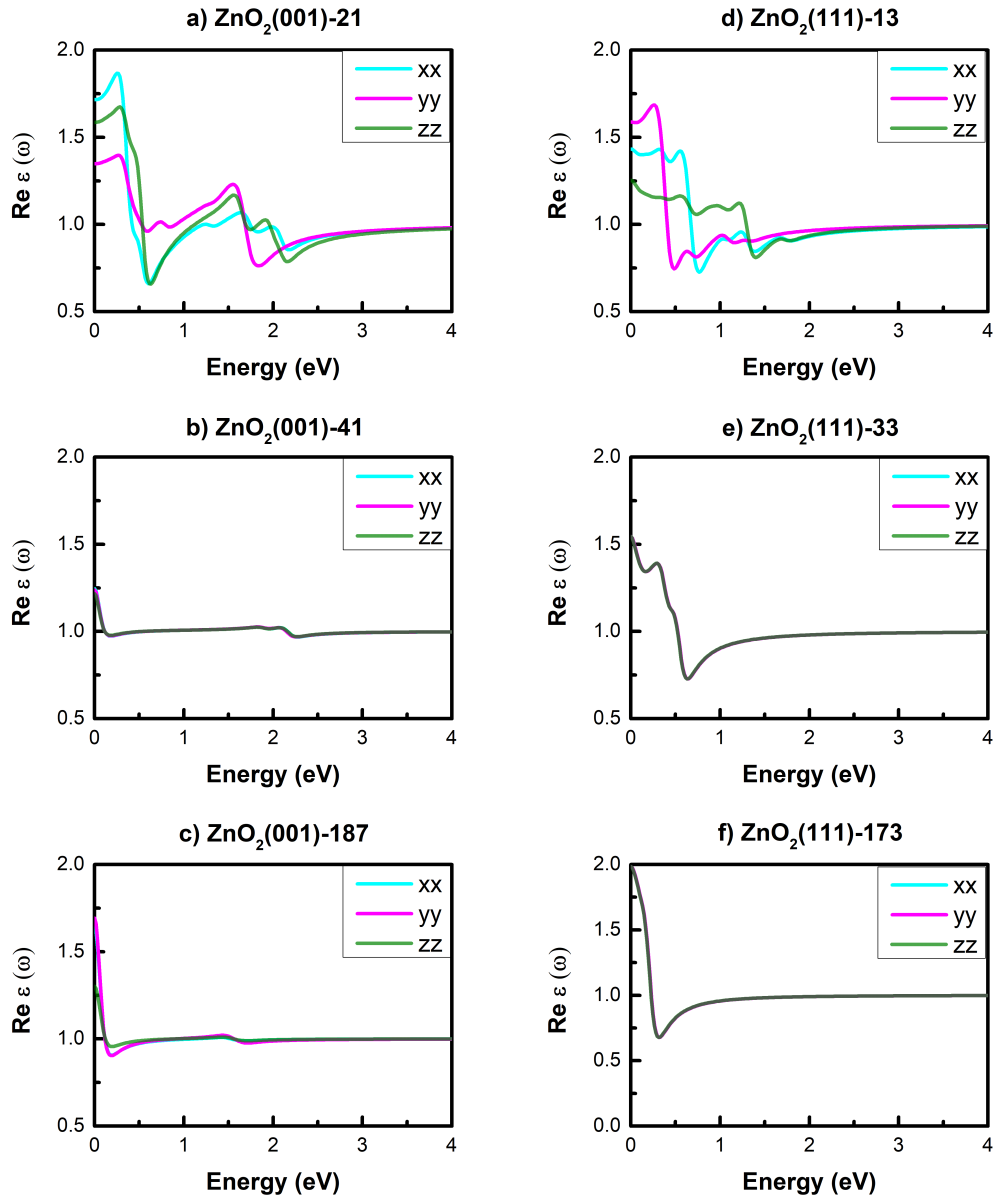


Fig. 4. Real part of the dielectric function of the ZnO₂ nanoparticles.

two smallest nanoparticles (ZnO₂(001)-21 and ZnO₂(111)-13 samples), we observe a noticeable spatial anisotropy of the real parts in the low energy range. Specifically, the maximum value of $\epsilon_1(\omega)$ for ZnO₂(001)-21 nanoparticles is approximately 2 in the x direction, while this value in the y direction is below 1.5. For ZnO₂(111)-13 nanoparticles, the dominant peak of $\epsilon_1(\omega)$ is about 1.75 in the y direction. In the visible light range (the incident photon energy is in between 1.65

eV and 3.2 eV), $\text{ZnO}_2(001)$ -21 nanoparticles show a better response of $\varepsilon_1(\omega)$ compared to that of $\text{ZnO}_2(111)$ -13.

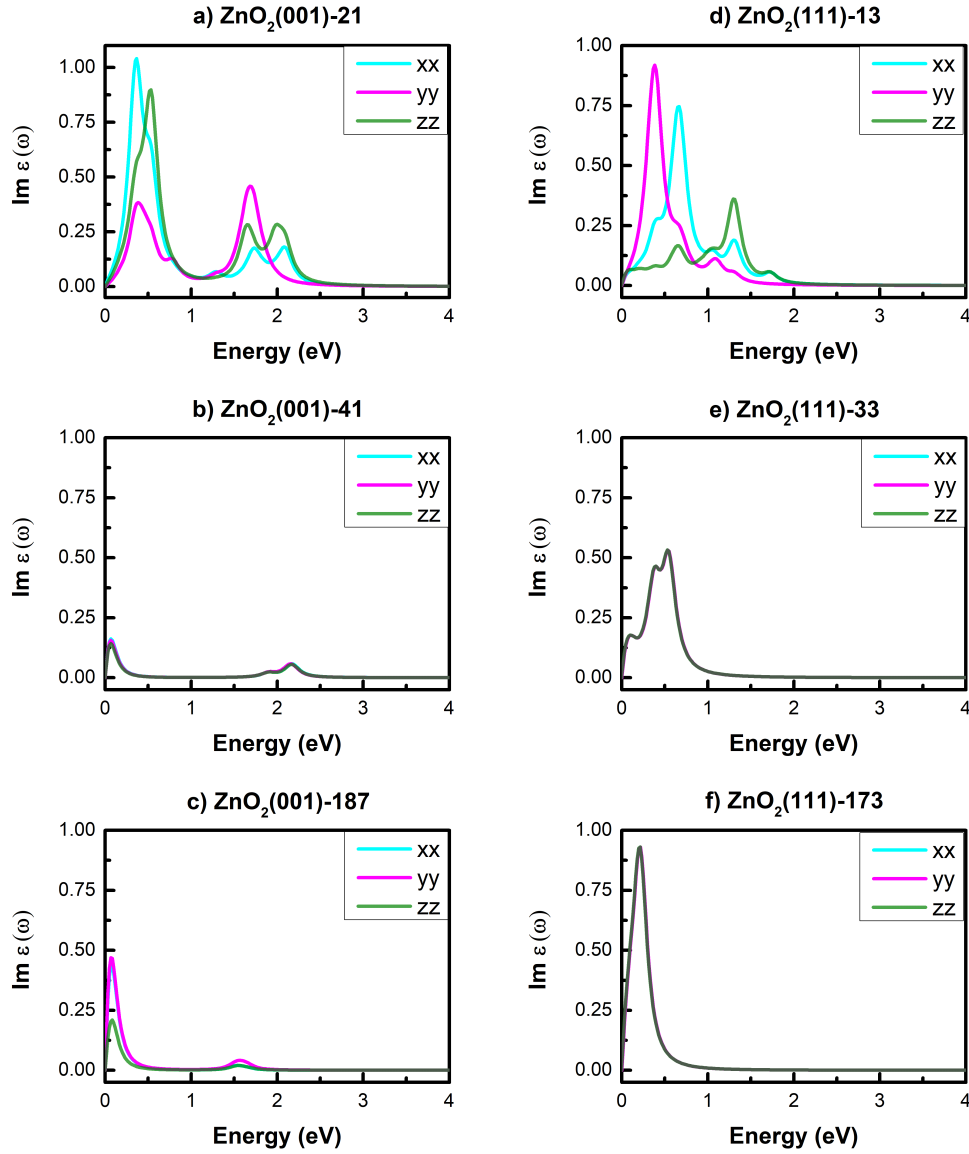


Fig. 5. Imaginary part of the dielectric function of the ZnO_2 nanoparticles.

We then show the imaginary parts of the dielectric function $\varepsilon_2(\omega)$ of all ZnO_2 samples in Fig. 5. These calculations share common optical features to the absorption coefficients presented in Fig. 6. Explicitly, the dominant absorption peaks of $\text{ZnO}_2(111)$ -33 and $\text{ZnO}_2(111)$ -173 appear in the energy range below 1.0 eV (mostly in the infrared region), whereas $\text{ZnO}_2(001)$ -21,

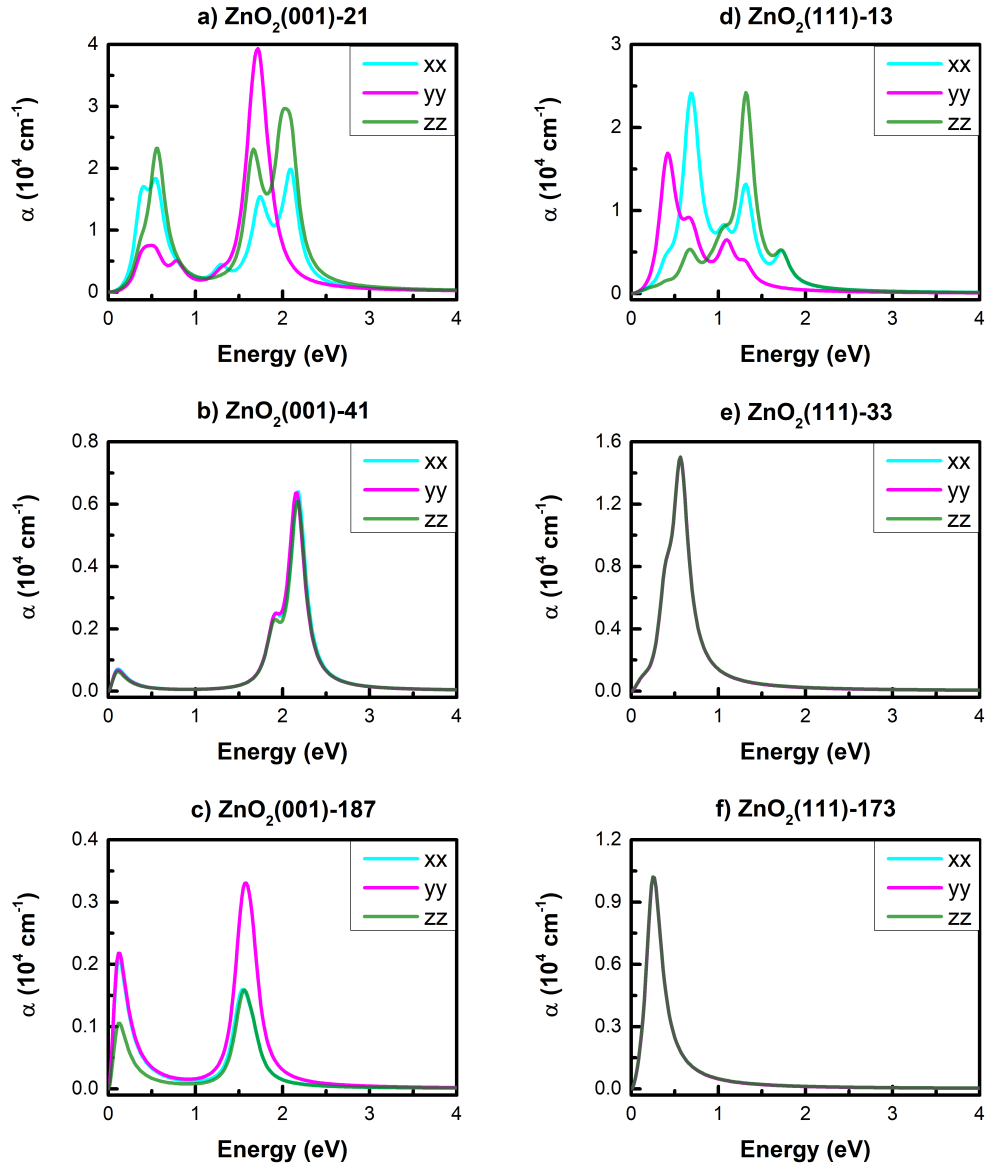


Fig. 6. Spatial anisotropic absorption coefficients of ZnO₂ nanoparticles.

ZnO₂(001)-41, ZnO₂(001)-187 and ZnO₂(111)-13 are more sensitive in the visible range. The imaginary part of the dielectric functions of ZnO₂(001)-21 and ZnO₂(111)-13 is highly asymmetry.

The absorption coefficient as a function of photon energy is calculated and shown in Fig. 6. To quantize the dominant peaks of these spectra with the electronic band structures, we sketch

these transition states in Fig. 7 and Fig. 8. It is well-known that the band-to-band absorption radiation can be generated by the photoexcitation of an electron being excited from a valence band to a conduction band. This transition can be direct or indirect. In the case of ZnO_2 nanoparticles, we associate the absorption coefficient spectra with direct transitions.

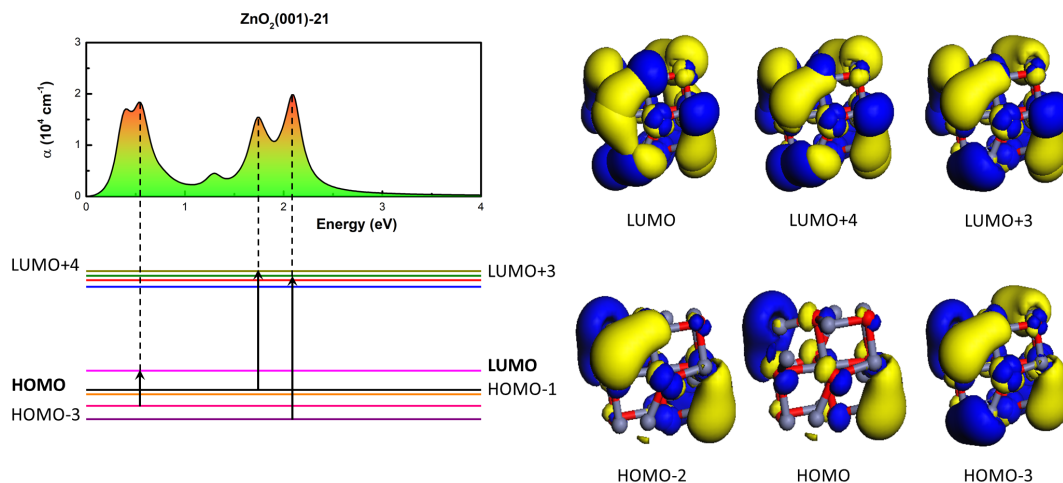


Fig. 7. Left: The electronic transition of states; Right: HOMO and LUMO (The isovalue set is 0.05) of several typical states for $\text{ZnO}_2(001)-21$.

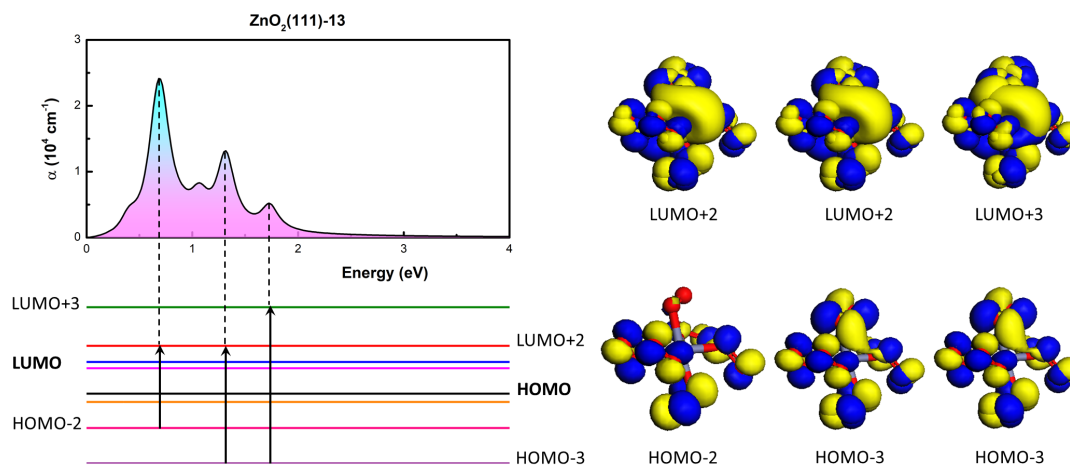


Fig. 8. Left: The electronic transition of states; Right: HOMO and LUMO (The isovalue set is 0.05) of several typical states for $\text{ZnO}_2(111)-13$.

For the case of ZnO₂ nanoparticles, except for ZnO₂(111)-33 and ZnO₂(111)-173 dominant peaks in the infrared region, the remaining nanoparticles exhibit dominant peaks in the visible light range, suggesting promising applications for early medical diagnosis devices. Furthermore, to elucidate the formation of these transition states, we present the band-to-band analyses and plot the highest occupied molecular orbital (HOMO) and the lowest unoccupied molecular orbital (LUMO) configurations of several typical states for two samples ZnO₂(001)-21 and ZnO₂(111)-13 in Fig. 7 and Fig. 8. For ZnO₂(001)-21 nanoparticles, there are the absorption peaks corresponding to electronic transitions: H-3 \rightarrow L+3, H \rightarrow L+4, H-2 \rightarrow L. The strongest peak occurs at an excitation energy of 2.1 eV corresponding to the electronic transition from H-3 \rightarrow L+3 in ZnO₂(001)-21. Also, the electronic transition from H-2 \rightarrow L matches the peak of the lowest intensity at 0.53 eV. In the case of ZnO₂(111)-13 nanoparticles, we obtain three peaks equivalent to the electronic transition from H-3 \rightarrow L+3, H-3 \rightarrow L+2 and H-2 \rightarrow L-1 at energy levels of 1.72, 1.3 and 0.7 eV, respectively. Thus, these results mark the novel optoelectronic properties of small or ultrasmall ZnO₂ nanoparticles. This demonstrates the richness of the optoelectronic properties of the small nanoparticles, which is not observed by ZnO₂ pure crystals [39].

IV. CONCLUSION

In summary, using DFT calculations, we have systematically investigated the opto-electronic properties of small nanoparticles, ZnO₂(001) and ZnO₂(111). By increasing the size and altering the surface orientations of the nanoparticles, we can tune the band gaps as well as the electronic properties of the ZnO₂ nanoparticles. Specifically, the two smallest nanoparticles exhibit semiconductor properties, while the rest of the investigated samples are electrically metal. Furthermore, albeit these nanoparticles are still small in size, we observe noticeable spatial anisotropies of absorption spectra indicating a reconciliation of the nanoparticle morphology and interfacial surface tensions. These results demonstrate the possibility of using ZnO₂ nanoparticles as an excellent candidate for developing the next generation of early diagnosis devices.

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