# THE BALANCING OF ROTATING MACHINERY AS NONLINEAR SYSTEM 

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#### Abstract

In this paper a method for dynamical balancing of rotating machines as a nonlinear system is proposed. After analysing and identifying the nonlinear system, a procedure of dynamical balance including measurement and processing of vibration signals, for calculating magnitude and location of imbalance mass is presented. An example of simulation and calculation is investigated for illustration of the method.


## 1. Introduction

One of the main causes of machinery vibration is imbalance of rotational parts. The dynamical balancing of rotating machines has been developing for new kinds of machines with high speed and improving balancing results more exactly and quickly. The system of machines with its rotational part is usually regarded as a linear system and there are some methods for balancing it. But sometimes it is difficult to find the location and magnitude of imbalance mass of the rotational part. Therefore, the balancing time is lasted so long and makes influence on production time. One of the reasons of this situation is the nonlinearity of the system. In this paper an effort is made for balancing a simple rotational part of machine with non-linear stiffness of the system.

## 2. The method for balancing rotational part of machine in linear

 caseThe brief of a dynamical balancing method for linear system is explained in the Fig. 1, where $m_{3}$ is rotational mass, $m_{4}$ and $m_{t}$ are eccentric and trial ones respectively. In this case the centrifugal force is proportional to according amplitude of vibration which can be measured. Therefore with three times of operating machine, it is able to find the location and magnitude of imbalance mass for balancing the rotational part of machine on the basis of triangle OAB (Fig.1), here $A_{0}$ is vibration amplitude for the case without trial mass and $A_{1}$, $A_{2}$ are amplitudes for the cases with trial mass.


Fig. 1

## 3. Balancing rotational part on non-linear system

### 3.1. Equation of motion

Let us investigate the system in Fig. 2
If the system has imbalance mass $m_{4}$ with distance $r$ from shaft line and the stiffness of the system is described by function $f(x)$ then the differential equation of motion in the direction of axis $X$ is governed as follows

$$
\begin{equation*}
m_{3} \ddot{x}+h_{1} \dot{x}+f(x)=m_{4} r \omega_{1}^{2} \cos \omega_{1} t \tag{3.1}
\end{equation*}
$$

Denoting $\quad \frac{h_{1}}{m_{3}}=2 h \omega_{0}$,

$$
\frac{f(x)}{m_{3}}=\omega_{0}^{2} x+g(x), \frac{m_{4}}{m_{3}} r=P
$$

we have


Fig. 2

$$
\begin{equation*}
\ddot{x}+2 h \omega_{0} \dot{x}+\omega_{0}^{2} x+g(x)=P \omega_{1}^{2} \cos \omega_{1} t \tag{3.2}
\end{equation*}
$$

In practice the function $f(x)$ and therefore the function $g(x)$ is symmetric for two sides of vibration with respect to axis $X$, therefore in mathematical expression the function $f(x)$ is odd, that means

$$
f(-x)=-f(x)
$$

Then the expand of $g(x)$ is of the form

$$
\begin{equation*}
g(x)=b_{3} x^{3}+b_{5} x^{5}+\ldots \tag{3.3}
\end{equation*}
$$

Since the vibration of the system is supposed to be small, then we restrict ourself by first term and the equation (2) becomes

$$
\begin{equation*}
\ddot{x}+2 h \omega_{0} \dot{x}+\omega_{0}^{2} x+b_{3} x^{3}=P \omega_{1}^{2} \cos \omega_{1} t \tag{3.4}
\end{equation*}
$$

where $h, \omega_{0}, \omega_{1}$ are given constants and $b_{3}$ is to be determined below.

### 3.2. Identification of the system by measurement and calculation

 of vibrationSuppose that $b_{3}$ is small and can be written

$$
\begin{equation*}
b_{3}=\varepsilon b, \tag{3.5}
\end{equation*}
$$

where $\varepsilon$ is a small positive parameter, we can write the equation (3.4) in the form

$$
\begin{equation*}
\ddot{x}+2 h \omega_{0} \dot{x}+\omega_{0}^{2} x=P \omega_{1}^{2} \cdot \cos \omega_{1} t-\varepsilon b x^{3} . \tag{3.6}
\end{equation*}
$$

On the basis of small parameter method [1] the solution of the equation (3.6) can be sought in the form

$$
\begin{equation*}
x(t)=x_{0}(t)+\varepsilon x_{1}(t)+\varepsilon^{2} x_{2}(t)+\ldots \tag{3.7}
\end{equation*}
$$

Substituting (3.7) into (3.6) then balancing the terms with the same order of $\varepsilon$ we obtain the following equations system:

$$
\begin{align*}
& \ddot{x}_{0}+2 h \omega_{0} \dot{x}_{0}+\omega_{0}^{2} x_{0}=P \omega_{1}^{2} \cos \omega_{1} t .  \tag{3.8}\\
& \ddot{x}_{1}+2 h \omega_{0} \dot{x}_{1}+\omega_{0}^{2} x_{1}=-b x_{0}^{3}  \tag{3.9}\\
& \ddot{x}_{2}+2 h \omega_{0} \dot{x}_{2}+\omega_{0}^{2} x_{2}=-3 b x_{0}^{2} x_{1} \tag{3.10}
\end{align*}
$$

We are interested in the stationary solutions of equations system (3.8)-(3.10). The stationary solution of Eq. (3.8) is described as follows

$$
\begin{equation*}
x_{0}(t)=\frac{P \omega_{1}^{2}}{D_{1}^{1 / 2}} \cos \phi_{0} \tag{3.11}
\end{equation*}
$$

where

$$
\begin{equation*}
D_{1}=\left(\omega_{0}^{2}-\omega_{1}^{2}\right)^{2}+4 h^{2} \omega_{0}^{2} \omega_{1}^{2}, \quad \phi_{0}=\omega_{1} t+\varphi_{0}, \quad \tan \varphi_{0}=\frac{2 h \omega_{0} \omega_{1}}{\omega_{1}^{2}-\omega_{0}^{2}} \tag{3.11}
\end{equation*}
$$

Equation (3.9) becomes

$$
\begin{align*}
& \ddot{x}_{1}+2 h \omega_{0} \dot{x}_{1}+\omega_{0}^{2} x_{1}=-\frac{b P^{3} \omega_{1}^{6}}{D_{1}^{3 / 2}} \cos ^{3} \phi_{0}  \tag{3.13}\\
& \cos ^{3} \phi_{0}=\frac{3}{4} \cos \phi_{0}+\frac{1}{4} \cos 3 \phi_{0} \tag{3.14}
\end{align*}
$$

and stationary solution of the Eq. (3.13) is

$$
\begin{equation*}
x_{1}(t)=-\frac{3 b P^{3} \omega_{1}^{6}}{4 D_{1}^{2}} \cos \left(\phi_{0}+\varphi_{11}\right)-\frac{b P^{3} \omega_{1}^{6}}{4 D_{1}^{3 / 2} D_{2}^{1 / 2}} \cos \left(3 \phi_{0}+\varphi_{13}\right), \tag{3.15}
\end{equation*}
$$

where

$$
D_{2}=\left(\omega_{0}^{2}-9 \omega_{1}^{2}\right)^{2}+36 h^{2} \omega_{0}^{2} \omega_{1}^{2}, \quad \tan \varphi_{11}=\tan \varphi_{0} \text { and } \tan \varphi_{13}=\frac{6 h \omega_{0} \omega_{1}}{9 \omega_{1}^{2}-\omega_{0}^{2}}
$$

It is similar for $x_{2}(t)$, which will be consisted of harmonics of first, third and fifth order of frequency $\omega_{1}$. We restrict ourselves by the first approximation. We have

$$
\begin{align*}
x(t) & =x_{0}(t)+\varepsilon x_{1}(t) \\
& =\frac{P \omega_{1}^{2}}{D_{1}^{1 / 2}} \cos \phi_{0}-\frac{3 P^{3} \omega_{1}^{6} b_{3}}{4 D_{1}^{2}} \cos \left(\phi_{0}+\varphi_{11}\right)-\frac{b_{3} P^{3} \omega_{1}^{6}}{4 D_{1}^{3 / 2} D_{2}^{1 / 2}} \cos \left(3 \phi_{0}+\varphi_{13}\right) \tag{3.16}
\end{align*}
$$

After some simple transformations the amplitudes of the first and the third order harmonics can be expressed by the following formulae.

$$
\begin{align*}
& A_{1}^{2}=\frac{P^{2} \omega_{1}^{4}}{D_{1}}+\frac{9 b_{3}^{2} P^{6} \omega_{1}^{12}}{16 D_{1}^{4}}-\frac{3 P^{4} \omega_{1}^{8} b_{3}\left(\omega_{1}^{2}-\omega_{0}^{2}\right)}{2 D_{1}^{3}}  \tag{3.17}\\
& A_{3}=\frac{b_{3} P^{3} \omega_{1}^{6}}{4 D_{1}^{3 / 2} D_{2}^{1 / 2}} \tag{3.18}
\end{align*}
$$

If it is possible to measure the vibration of the system along the axis $X$ and describe it in frequency domain, then it is able to obtain $A_{1}$ for frequency $\omega_{1}$ and $A_{3}$ for frequency $3 \omega_{1}$ from measurement data.

Using two equations (3.17) and (3.18) we have the following equation for $q=P \omega_{1}^{2}:$

$$
\begin{equation*}
q^{2}-6 A_{3}\left(\omega_{1}^{2}-\omega_{0}^{2}\right) \frac{D_{2}^{1 / 2}}{D_{1}^{1 / 2}} q+9 A_{3}^{2} D_{2}-A_{1}^{2} D_{1}=0 \tag{3.19}
\end{equation*}
$$

From here we obtain

$$
\begin{equation*}
q=P \omega_{1}^{2}=3 A_{3}\left(\omega_{1}^{2}-\omega_{0}^{2}\right) \frac{D_{2}^{1 / 2}}{D_{1}^{1 / 2}} \pm \sqrt{\Delta} \tag{3.20}
\end{equation*}
$$

where $\Delta=9 A_{3}^{2}\left(\omega_{1}^{2}-\omega_{2}^{2}\right)^{2} \frac{D_{2}}{D_{1}}-9 A_{3}^{2} D_{2}+A_{1}^{2} D_{1}$ and from (3.18) we have the value $b_{3}$.

### 3.3. Determination of balancing parameters

Adding a trial mass $m_{t}$ alternatively at two radial symmetric places of rotational part, measuring vibrations for every case and using similar formula as (3.20) we can calculate the values $P_{1}$ and $P_{2}$ which are proportional to centrifugal forces.

Draw the triangle, which is similar to triangle OAB on Fig. 1 and allows us to determine the angle $\varphi$, that means, the location of imbalance mass $m_{4}$. The magnitude of $m_{4}$ is calculated by formula

$$
\begin{equation*}
m_{4}=\frac{m_{3} P}{r} \tag{3.21}
\end{equation*}
$$

We can check it by computing $P_{t}$ from the expression

$$
\begin{align*}
P_{t}^{2} & =\frac{P_{1}^{2}+P_{2}^{2}-2 P^{2}}{2}  \tag{3.22}\\
m_{4} & =\frac{m_{t} P}{P_{t}} \tag{3.23}
\end{align*}
$$

and the angle $\varphi$ can be calculated by the formula

$$
\begin{equation*}
\varphi=\arccos \left[\frac{P_{1}^{2}-P_{2}^{2}}{4 P P_{t}}\right] \tag{3.24}
\end{equation*}
$$

### 3.4. Procedure for balancing rotational part in non-linear system

- Measure three times vibrations of bearing on imbalance system, determine
amplitudes at frequency $\omega_{1}$ and $3 \omega_{1}$ on graphics of vibrations in frequency domain for initial imbalance mass and two times of attachment of the trial mass.
- Determine $P, P_{1}, P_{2}$ for three times of measurement by formula similar to (3.20)
- Calculate $P_{t}$ and $\varphi$ from formulas (3.22) and (3.24). Two balancing parameters $m_{4}$ and $\varphi$ are determined.
- Attach the balance mass $m_{4}$ (calculated by (3.21)) in the place of radial symmetry with respect to position found by angle $\varphi$, the rotational part will be balanced.


## 4. Results of simulation and calculation

- The equation (3.4) is solved directly by computer programme with given parameters including $b_{3}$ and $P$, and the solution (vibration) and its spectrum can be obtained. We can determine two values $A_{1}$ and $A_{3}$ (Figs 3-5).
- Using input parameters: $\omega_{1}, \omega_{0}, h, A_{1}, A_{3}$.

That means, it is not necessary to know nonlinear coefficient $b_{3}$ and intense of force $P$, but we can get these values from equations $(3,18)$ and $(3.20)$.


Fig. 3. Vibration and spectrum
$h_{1}=0.029811, m_{4}=1 ., f=4.74185, f_{1}=8 ., \beta=0.5, r=0.6$
Ordinate level: $\mathrm{E}=56.4559, \mathrm{D}=42.34192, \mathrm{C}=28.22795, \mathrm{~B}=14.11397, \mathrm{~A}=0.0$


Fig. 4. Vibration and spectrum
$h_{1}=0.02981, m_{4}=1.5, f=4.74173, f_{1}=8 ., \beta=0.5, r=0.6$
Ordinate level : $\mathrm{E}=60.19252, \mathrm{D}=45.14439, \mathrm{C}=30.09626, \mathrm{~B}=15.04813, \mathrm{~A}=0.0$


Fig. 5. Vibration and spectrum
$h_{1}=0.029809, m_{4}=2 ., f=4.74162, f_{1}=8 ., \beta=0.5, r=0.6$
Ordinate level : $\mathrm{E}=63.34972, \mathrm{D}=47.51229, \mathrm{C}=31.67486, \mathrm{~B}=15.83743, \mathrm{~A}=0.0$

- Comparison of the values $P_{g}$ (calculated by formula $P_{g}=\frac{m_{4}}{m_{3}} r$ ) with computational results $P_{c}$ (by (3.20)) is given in the Table 1 with amplify coefficient $10^{4}$, in which $\omega_{1}=2 \pi f_{1}, \omega_{0}=2 \pi f$ and $f_{1}=8, f$ and $h$ are calculated from
simulation model, $m_{3}=20,000 \mathrm{~kg}$. The values in two final columns are closed each other, that shows the exactitude of the method.

Table 1

| $m_{4}$ | $f$ | $h$ |  | $A_{1}$ |  | $A_{3}$ |  | $P_{g}$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  | $P_{c}$ |  |  |  |
| 2.50 | 4.74150 | 0.029809 | 66.09504 |  | 4.55542 |  | 75.5727 | 74.8783 |
| 2.25 | 4.74156 | 0.029809 | 64.76664 |  | 4.22026 |  | 67.6920 | 67.3913 |
| 2.00 | 4.74162 | 0.029809 | 63.34972 |  | 3.87782 | 59.9041 | 59.6823 |  |
| 1.75 | 4.74167 | 0.029810 | 61.83076 | 3.54005 | 52.4168 | 51.8845 |  |  |
| 1.50 | 4.74173 | 0.029810 | 60.19252 | 3.20013 | 44.9292 | 44.0758 |  |  |
| 1.25 | 4.74179 | 0.029811 | 58.41179 |  | 2.85299 | 37.4415 | 36.1874 |  |
| 1.00 | 4.74185 | 0.029811 | 56.45590 | 2.51983 | 29.9536 | 28.7850 |  |  |
| 0.75 | 4.74191 | 0.029811 | 54.27776 | 2.19330 | 22.4855 | 21.7034 |  |  |
| 0.50 | 4.74196 | 0.029812 | 51.80022 | 1.84879 | 14.9772 | 14.3523 |  |  |

## 5. Conclusions

The method presented in this paper is rather simple for practical application, after computing $P_{c}$ (final column of Table 1) the results for linear system could be applied and the balancing parameters are found out.

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## REFERENCES

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## CÂN BÅNG MÁY QUAY TRONG TRƯỜNG HỢP PHI TUYẾN

Bài này đề xuất một phương pháp cân bằng động máy quay trong trường hợp phi tuyến. Sau khi phân tích và nhận dạng hệ phi tuyến, một quy trình cân bẵng động bao gồm đo dao động, xử lý tín hiệu và dùng các tham số để tính toán độ lớn và vị trí khối lượng mất cân bằng đã được trình bày. Một ví dụ về mô phỏng và tính toán đã được thực hiện để minh họa cho phương pháp.

