# TWO-DIMENSIONAL OPTIMIZATION PROBLEM OF PLANT LOCATION 

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#### Abstract

In this paper, the following matters are presented: the adjoint problem of the two-dimensional matter propagation problem; the algorithm for determination of a domain in which a plant can be located so that the values of the pollution-level reflecting functional does not exceed a given value at considered sensitive areas; application of this algorithm for numerical experiments to a typical problem.


1. Equation of the suspended matter propagation and its adjoint equation (see [1])

The equation of the suspended matter propagation, i.e. the matter transport and diffusion equation in the horizontal 2D case has the following form:

$$
\begin{equation*}
\frac{\partial C}{\partial t}+u \frac{\partial C}{\partial x}+v \frac{\partial C}{\partial y}+\sigma C=f+\gamma \Delta C, \quad(x, y) \in G, 0 \leq t \leq T \tag{1.1}
\end{equation*}
$$

with the initial and boundary conditions:

$$
\begin{equation*}
\left.C\right|_{t=0}=C^{0},\left.\quad C\right|_{\Gamma^{-}}=\varphi,\left.\quad \frac{\partial C}{\partial n}\right|_{\Gamma^{+}}=0 \tag{1.2}
\end{equation*}
$$

where $x, y, t$ are the space and time variables; $C$ is the matter concentration; $\sigma$ is the decay coefficient; $f$ is the source intensity; $\gamma$ is the diffusion coefficient; $u, v$ are respectively velocity components in the $x$ and $y$ directions, and satisfy the following equation:

$$
\begin{equation*}
\frac{\partial u}{\partial x}+\frac{\partial v}{\partial y}=0 \tag{1.3}
\end{equation*}
$$

$\Gamma=\Gamma^{+}+\Gamma^{-}$with $\Gamma^{+}$is the boundary part, at which $u_{n} \geq 0 ; \Gamma^{-}$is the boundary part, at which $u_{n}<0 ; \Delta=\frac{\partial^{2}}{\partial x^{2}}+\frac{\partial^{2}}{\partial y^{2}}$-Laplace operator; $u_{n}$ is the projection of the velocity on the external normal vector $\vec{n}$.

Using (1.3), the equation (1.1) can be rewritten as follows:

$$
\begin{equation*}
\frac{\partial C}{\partial t}+\frac{\partial u C}{\partial x}+\frac{\partial v C}{\partial y}+\sigma C=f+\gamma \Delta C \tag{1.4}
\end{equation*}
$$

Solution of the equation (1.1) may be determined under the form: $C=C_{1}+C_{2}$ where, $C_{1}$ and $C_{2}$ are the solutions of the following equation:

$$
\begin{equation*}
\frac{\partial C_{1}}{\partial t}+u \frac{\partial C_{1}}{\partial x}+v \frac{\partial C_{1}}{\partial y}+\sigma C_{1}=\gamma \Delta C_{1} \tag{1.5}
\end{equation*}
$$

with the initial and boundary conditions:

$$
\left.C_{1}\right|_{t=0}=C^{0},\left.\quad C_{1}\right|_{\Gamma^{-}}=\varphi,\left.\quad \frac{\partial C_{1}}{\partial n}\right|_{\Gamma^{+}}=0
$$

and

$$
\begin{equation*}
\frac{\partial C_{2}}{\partial t}+\frac{\partial u C_{2}}{\partial x}+\frac{\partial v C_{2}}{\partial y}+\sigma C_{2}=\gamma \Delta C_{2}+f \tag{1.6}
\end{equation*}
$$

with the initial and boundary conditions:

$$
\begin{equation*}
\left.C_{2}\right|_{t=0}=0,\left.\quad C_{2}\right|_{\Gamma^{-}}=0,\left.\quad \frac{\partial C_{2}}{\partial n}\right|_{\Gamma^{+}}=0 . \tag{1.7}
\end{equation*}
$$

We now establish the adjoint equation of the equation (1.6). By multiplying both sides of the equation (1.6) by a some function $C_{2}^{*}$ and integrating the equation obtained on the area $G \times[0, T]$, we get:

$$
\begin{align*}
\int_{0}^{T} d t \int_{G} C_{2}^{*} & \frac{\partial C_{2}}{\partial t} d G+\int_{0}^{T} d t \int_{G} C_{2}^{*} \operatorname{div}\left(\vec{u} C_{2}\right) d G \\
& +\int_{0}^{T} d t \int_{G} \sigma C_{2} C_{2}^{*} d G-\gamma \int_{0}^{T} d t \int_{G} C_{2}^{*} \Delta C_{2} d G=\int_{0}^{T} d t \int_{G} C_{2}^{*} f d G \tag{1.8}
\end{align*}
$$

Let $\gamma=$ const, using the partial integration technique, the Green formula and the condition (1.3), we have:

$$
\begin{aligned}
& \int_{0}^{T} d t \int_{G} C_{2}^{*} \frac{\partial C_{2}}{\partial t} d G=\left.\int_{G} C_{2}^{*} C_{2}\right|_{0} ^{T} d G-\int_{0}^{T} d t \int_{G} C_{2} \frac{\partial C_{2}^{*}}{\partial t} d G \\
& \int_{0}^{t} d t \int_{G} C_{2}^{*} \operatorname{div}\left(\vec{u} C_{2}\right) d G=\int_{0}^{T} d t \int_{\Gamma}^{T} u_{n} C_{2}^{*} C_{2} d \Gamma-\int_{0}^{T} d t \int_{G} C_{2} \operatorname{div}\left(\vec{u} C_{2}^{*}\right) d G \\
& \gamma \int_{0}^{T} d t \int_{G} C_{2}^{*} \Delta C_{2} d G=\gamma \int_{0}^{T} d t \int_{\Gamma}\left(C_{2}^{*} \frac{\partial C_{2}}{\partial n}-C_{2} \frac{\partial C_{2}^{*}}{\partial n}\right) d \Gamma+\gamma \int_{0}^{T} d t \int_{G} C_{2} \Delta C_{2}^{*} d G
\end{aligned}
$$

Putting these expressions into (1.8), one deduces:

$$
\begin{align*}
& \int_{0}^{T} d t \int_{G} C_{2}\left(-\frac{\partial C_{2}^{*}}{\partial t}-\operatorname{div}\left(\vec{u} C_{2}^{*}\right)+\sigma C_{2}^{*}-\gamma \Delta C_{2}^{*}\right) d G= \\
& =\int_{0}^{T} d t \int_{G} C_{2}^{*} f d G-\left.\int_{G} C_{2}^{*} C_{2}\right|_{t=T} d G+\left.\int_{G} C_{2}^{*} C_{2}\right|_{t=0} d G-\int_{0}^{T} d t \int_{\Gamma} u_{n} C_{2} C_{2}^{*} d \Gamma \\
& \quad+\gamma \int_{0}^{T} d t \int_{\Gamma}\left(C_{2}^{*} \frac{\partial C_{2}}{\partial n}-C_{2} \frac{\partial C_{2}^{*}}{\partial n}\right) d \Gamma \tag{1.9}
\end{align*}
$$

Let the function $C_{2}^{*}$ satisfy the following equation:

$$
\begin{equation*}
-\frac{\partial C_{2}^{*}}{\partial t}-\operatorname{div}\left(\vec{u} C_{2}^{*}\right)+\sigma C_{2}^{*}-\gamma \Delta C_{2}^{*}=p \tag{1.10}
\end{equation*}
$$

From the initial and the boundary conditions (1.7), one yields;

$$
\begin{aligned}
& \left.\int_{G} C_{2} C_{2}^{*}\right|_{t=0} d G=0, \\
& \int_{0}^{T} d t \int_{\Gamma} u_{n} C_{2} C_{2}^{*} d \Gamma=\int_{0}^{T} d t \int_{\Gamma^{+}} u_{n} C_{2} C_{2}^{*} d \Gamma \\
& \gamma \int_{0}^{T} d t \int_{\Gamma}\left(C_{2}^{*} \frac{\partial C_{2}}{\partial n}-C_{2} \frac{\partial C_{2}^{*}}{\partial n}\right) d \Gamma=\gamma \int_{0}^{T} d t \int_{\Gamma^{+}}\left(-C_{2} \frac{\partial C_{2}^{*}}{\partial n}\right) d \Gamma+\gamma \int_{0}^{T} d t \int_{\Gamma^{-}} C_{2}^{*} \frac{\partial C_{2}}{\partial n} d \Gamma .
\end{aligned}
$$

From the above expressions and (1.10), the equation (1.9) can be rewritten under the form:

$$
\begin{align*}
\int_{0}^{T} d t \int_{G} p C_{2} d G= & \int_{0}^{T} d t \int_{G} f C_{2}^{*} d G+\left.\int_{G} C_{2} C_{2}^{*}\right|_{t=T} d G \\
& +\gamma \int_{0}^{T} d t \int_{\Gamma^{-}} C_{2}^{*} \frac{\partial C_{2}}{\partial n} d \Gamma-\int_{0}^{T} d t \int_{\Gamma^{+}} C_{2}\left(\gamma \frac{\partial C_{2}^{*}}{\partial n}+u_{n} C_{2}^{*}\right) d \Gamma \tag{1.11}
\end{align*}
$$

Let the initial and boundary conditions of the equation (1.10) be chosen as follows:

$$
\begin{equation*}
\left.C_{2}^{*}\right|_{t=T}=0,\left.\quad C_{2}^{*}\right|_{\Gamma^{-}}=0,\left.\quad\left(\gamma \frac{\partial C_{2}^{*}}{\partial n}+u_{n} C_{2}^{*}\right)\right|_{\Gamma^{+}}=0 \tag{1.12}
\end{equation*}
$$

Then, from (1.11) and (1.12) we get the dual form:

$$
\begin{equation*}
\int_{0}^{T} d t \int_{G} p C_{2} d G=\int_{0}^{T} d t \int_{G} f C_{2}^{*} d G \tag{1.13}
\end{equation*}
$$

It is easy to verify that the problem $(1.10),(1.12)$ is the adjoint problem of the (1.6), (1.7). Indeed, with the notation:

$$
A=\frac{\partial}{\partial t}+\frac{\partial u \cdot}{\partial x}+\frac{\partial v}{\partial y}+\sigma-\gamma \Delta \quad A^{*}=-\frac{\partial}{\partial t}-\frac{\partial u \cdot}{\partial x}-\frac{\partial v}{\partial y}+\sigma-\gamma \Delta
$$

we have:

$$
\begin{aligned}
& A C_{2}=f, \quad A^{*} C_{2}^{*}=p \\
& \left(A C_{2}, C_{2}^{*}\right)=\left(f, C_{2}^{*}\right)=\int_{0}^{T} d t \int_{G} f C_{2}^{*} d G=\int_{0}^{T} d t \int_{G} p C_{2} d G=\left(C_{2}, p\right)=\left(C_{2}, A^{*} C_{2}^{*}\right)
\end{aligned}
$$

Use of the variable transformation $t_{1}=T-t$, the equation (1.10) becomes:

$$
\begin{align*}
& \frac{\partial C_{2}^{*}}{\partial t_{1}}-\operatorname{div}\left(\vec{u} C_{2}^{*}\right)+\sigma C_{2}^{*}-\gamma \Delta C_{2}^{*}=p  \tag{1.14}\\
& \left.C_{2}^{*}\right|_{t_{1}=0}=0,\left.\quad C_{2}^{*}\right|_{\Gamma^{-}}=0,\left.\quad\left(\gamma \frac{\partial C_{2}^{*}}{\partial n}+u_{n} C_{2}^{*}\right)\right|_{\Gamma^{+}}=0
\end{align*}
$$

For simplicity, by using (1.3), we obtain an another form of the adjoint equation (1.14):

$$
\begin{equation*}
\frac{\partial C_{2}^{*}}{\partial t_{1}}-u \frac{\partial C_{2}^{*}}{\partial x}-v \frac{\partial C_{2}^{*}}{\partial y}+\sigma C_{2}^{*}-\gamma \Delta C_{2}^{*}=p \tag{1.15}
\end{equation*}
$$

## 2. Pollution-level reflecting functionals (see [1])

Assume that the suspended matter concentration $C$ is calculated from the equation (1.1). We consider the following functionals:
a. The time-averaged amount of the matter concentration $C$ on a sensitive area $G_{k} \subset G$ for the period $T: J_{k}^{A}=\frac{1}{T} \int_{0}^{T} d t \int_{G_{k}} C d G$.
b. The total amount of settling matter in the same area $G_{k} \subset G: J_{k}^{B}=\int_{0}^{T} d t \int_{G_{k}} a C d G$, where, the constant a represents portion of matter which settles down, that are mainly the heavy matters and partly the suspended matters settling down by downward diffusion.
c. Generalized functional:

$$
J_{k}=\int_{0}^{T} d t \int_{G_{k}} p C d G \quad \text { where } \quad p= \begin{cases}\frac{1}{T}+a, & (x, y) \in G_{k}  \tag{2.1}\\ 0, & (x, y) \notin G_{k}\end{cases}
$$

and $p$ is a function referring to the economic, sanitary, ecological, health standards and so on.
d. Global functional:

$$
Y_{p}=\int_{0}^{T} d t \int_{G} p C d G \text { where, } \quad p= \begin{cases}\frac{1}{T}+G_{k}, & (x, y) \in G_{k}, k=1,2, \ldots, m \\ 0, & (x, y) \notin \bigcup_{k=1}^{m} G_{k} .\end{cases}
$$

## 3. Optimization problem of plant location (see [1])

Let $G_{k}(k=1,2, \ldots, m)$ be considered areas, recreation zones or other environmentally sensitive areas on the region $G$. Our problem is to determine the domain $\Omega_{k} \subset G$ so that the pollution matter from a plant located in this domain $\Omega_{k}$ satisfies the following condition for the sensitive area $G_{k}$ :

$$
Y_{k}=\int_{0}^{T} d t \int_{G_{k}} p C d G \leq \bar{c}_{k}, \quad \text { where } \quad p= \begin{cases}\frac{1}{T}+a_{k}, & (x, y) \in G_{k}  \tag{3.1}\\ 0, & (x, y) \notin G_{k}\end{cases}
$$

and $\bar{c}_{k}$ is a given figure.
If the determination of domain $\Omega_{k}$ is impossible on the $G$, the reduction of rate of the pollution emission $Q$, will make the determination of the plant location possible.

Assume that on the region $G$ there are $m$ sensitive areas $G_{k}(k=1, \ldots, m)$ and the source of matter emission is located at a point $r_{0}=\left(x_{0}, y_{0}\right)$. Then, the source intensity can be described by the function: $f(x, y)=Q \delta\left(r-r_{0}\right), \quad Q=$ const where, $\delta(r)=\left\{\begin{array}{ll}\infty, & r=r_{0} \\ 0, & r \neq r_{0}\end{array}\right.$ is Dirac function, and from (1.1), we get:

$$
\frac{\partial C}{\partial t}+u \frac{\partial C}{\partial x}+v \frac{\partial C}{\partial y}+\sigma C=Q \delta\left(r-r_{0}\right)+\gamma \Delta C
$$

with the conditions: $\left.C\right|_{t=0}=C^{0},\left.\quad C\right|_{\Gamma^{-}}=\varphi,\left.\quad \frac{\partial C}{\partial n}\right|_{\Gamma^{+}}=0$.
In order to determine the domain $\Omega$, in which the plant can be located so that in all sensitive areas $G_{k}$, the generalized functional $Y_{k}$ satisfies the condition (3.1), we do as follows:
a. Calculation of concentration $C$ from the equation (1.5):

$$
\frac{\partial C}{\partial t}+u \frac{\partial C}{\partial x}+v \frac{\partial C}{\partial y}+\sigma C=\gamma \Delta C
$$

with the initial and boundary conditions $\left.C\right|_{t=0}=C^{0},\left.\quad C\right|_{\Gamma^{-}}=\varphi,\left.\quad \frac{\partial C}{\partial n}\right|_{\Gamma^{+}}=0$ and generalized functional (2.1): $J_{k}=\int_{0}^{T} d t \int_{G_{k}} p C d G=\widetilde{c}_{k}$.
b. Solving $m$ adjoint equations:

$$
\frac{\partial C_{k}^{*}}{\partial t_{1}}-u \frac{\partial C_{k}^{*}}{\partial x}-v \frac{\partial C_{k}^{*}}{\partial y}+\sigma C_{k}^{*}-\gamma \Delta C_{k}^{*}=p_{k}
$$

where, $p_{k}=\left\{\begin{array}{ll}\frac{1}{T}+a_{k}, & (x, y) \in G_{k} \\ 0, & (x, y) \notin G_{k}\end{array}\right.$ with the conditions:

$$
\left.C_{k}^{*}\right|_{t_{1}=0}=0,\left.\quad C_{k}^{*}\right|_{\Gamma^{-}}=0,\left.\quad\left(v \frac{\partial C_{k}^{*}}{\partial n}+u_{n} C_{k}^{*}\right)\right|_{\Gamma^{+}}=0
$$

we obtain the solutions $C_{k}^{*}(k=1,2, \ldots, m)$.
From the dual form (1.13), we get:

$$
\begin{aligned}
Y_{k}^{*} & =\int_{0}^{T} d t \int_{G} p_{k} C d G=\int_{0}^{T} d t \int_{G} Q \delta\left(r-r_{0}\right) C_{k}^{*} d G \\
& =\int_{0}^{T} Q C_{k}^{*}\left(r_{0}, t\right) d t=\int_{0}^{T} Q C_{k}^{*}\left(r_{0}, T-t_{1}\right) d t_{1}
\end{aligned}
$$

which must satisfy the condition: $Y_{k}^{*} \leq \bar{c}_{k}-\widetilde{c}_{k}=\overline{\bar{c}}_{k}$.
Now we consider the function: $Y_{k}^{*}(r)=Q \int_{0}^{T} C_{k}^{*}(r, t) d t$ and draw the iso-grams of $Y_{k}^{*}(r)=$ const. Then, $\Omega_{k}$ in which the functional $Y_{k}^{*}(r) \leq \overline{\bar{c}}_{k}$ are found out. If there is perchance no area $\Omega_{k}$ inside $G$, it may be re-established anyway by reducing the discharge intensity $Q$.
c. Overlaying all the areas $\Omega_{k}(k=1, \ldots, m)$, we obtain the domain $\Omega$, $\left(\Omega=\bigcap_{k=1}^{m} \Omega_{k}\right)$. $\Omega$ will be the domain in which the plant can be located so that pollution standards will be met in all the areas $G_{k} \subset G,(k=1,2, \ldots, m)$.

## 4. Algorithm (see [2]-[4])

The equation (1.5) and the adjoint equation (1.15) may be rewritten in a common form:

$$
\begin{equation*}
\frac{\partial C}{\partial t}+\Lambda C=f \tag{4.1}
\end{equation*}
$$

where, $\Lambda=\Lambda_{1}+\Lambda_{2}, \Lambda_{1}= \pm u \frac{\partial}{\partial x}-\gamma \frac{\partial^{2}}{\partial x^{2}}+\frac{\sigma}{2}, \Lambda_{2}= \pm v \frac{\partial}{\partial y}-\gamma \frac{\partial^{2}}{\partial y^{2}}+\frac{\sigma}{2}$.
Equation (4.1) is solved by the method of the directional decomposition (splitting method):

$$
\begin{align*}
& \frac{C^{k+1}-C^{k}}{d t}+\Lambda\left[\theta C^{k+1}+(1-\theta) C^{k}\right]=f^{k+1} \\
\text { or } & (I+d t \theta \Lambda) C^{k+1}=[I-d t(1-\theta) \Lambda] C^{k}+d t f^{k+1} \tag{4.2}
\end{align*}
$$

where $0 \leq \theta \leq 1, I$ is the unique operator.
Using approximation:

$$
\left[I+d t \theta\left(\Lambda_{1}+\Lambda_{2}\right)\right]=\left(I+d t \theta \Lambda_{1}\right)\left(I+d t \theta \Lambda_{2}\right)+0\left(d t^{2}\right)
$$

from (4.2), one deduces:

$$
\left(I+d t \theta \Lambda_{1}\right)\left(I+d t \theta \Lambda_{2}\right) C^{k+1}=d t f^{k+1}+[I-d t(1-\theta) \Lambda] C^{k}
$$

The computational process contains two steps:

$$
\begin{align*}
\left(I+d t \theta \Lambda_{1}\right) C^{k+1 / 2} & =[I-d t(1-\theta) \Lambda] C^{k}+d t f^{k+1}  \tag{4.3}\\
\left(I+d t \theta \Lambda_{2}\right) C^{k+1} & =C^{k+1 / 2} \tag{4.4}
\end{align*}
$$

a. Discretizing the equation (4.3) by an implicit finite difference scheme:

$$
\begin{aligned}
& \Lambda_{1} C^{k+1 / 2}=\frac{( \pm u+|u|)_{m, n}^{k+1 / 2}}{2} \frac{\left(C_{m, n}^{k+1 / 2}-C_{m-1, n}^{k+1 / 2}\right)}{\Delta x} \\
& +\frac{( \pm u-|u|)_{m, n}^{k+1 / 2}}{2} \frac{\left(C_{m+1, n}^{k+1 / 2}-C_{m, n}^{k+1 / 2}\right)}{\Delta x}-\gamma \frac{\left(C_{m+1, n}^{k+1 / 2}-2 C_{m, n}^{k+1 / 2}+C_{m-1, n}^{k+1 / 2}\right)}{\Delta x^{2}}+\frac{\sigma}{2} \\
& \Lambda C^{k}= \pm u_{m, n}^{k+1 / 2} \frac{\left(C_{m+1, n}^{k}-C_{m-1, n}^{k}\right)}{2 \Delta x}-\gamma \frac{C_{m+1, n}^{k}-2 C_{m, n}^{k}+C_{m-1, n}^{k}}{\Delta x^{2}} \\
& \quad \pm v_{m, n}^{k+1 / 2} \frac{\left(C_{m, n}^{k+1}-C_{m, n-1}^{k}\right)}{2 \Delta y}-\gamma \frac{C_{m, n+1}^{k}-2 C_{m, n}^{k}+C_{m, n-1}^{k}}{\Delta y^{2}}+\sigma
\end{aligned}
$$

we obtain:

$$
\begin{equation*}
a_{m} C_{m+1, n}^{k+1 / 2}+b_{m} C_{m, n}^{k+1 / 2}+c_{m} C_{m-1, n}^{k+1 / 2}=d_{m} \tag{4.5}
\end{equation*}
$$

where,

$$
\begin{aligned}
& a_{m}=\frac{( \pm u-|u|)_{m, n}^{k+1 / 2} \theta d t}{2 \Delta x}-\frac{\gamma \theta d t}{(\Delta x)^{2}}, \quad b_{m}=1+\frac{\theta|u|_{m, n}^{k+1 / 2} d t}{\Delta x}+2 \frac{\gamma \theta d t}{(\Delta x)^{2}}+\frac{\sigma d t}{2} \\
& c_{m}=-\frac{( \pm u+|u|)_{m, n}^{k+1 / 2} \theta d t}{2 \Delta x}-\frac{\gamma \theta d t}{(\Delta x)^{2}}, \quad d_{m}=d t f_{m, n}^{k+1}+[I-d t(1-\theta) \Lambda] C_{m, n}^{k}
\end{aligned}
$$

It is easy to verify that: $b_{m}>0, a_{m}<0, c_{m}<0$ and $\left|b_{m}\right| \geq\left|a_{m}\right|+\left|c_{m}\right|+\delta$, $\delta>0$.

So, the linear equation system (4.5) has the unique solution and the computational error of the following double sweep method

$$
\begin{equation*}
C_{m, n}^{k+1}=L_{m} C_{m+1, n}^{k+1}+K_{m} \tag{4,6}
\end{equation*}
$$

where, $L_{m}=\frac{-a_{m}}{b_{m}+c_{m} L_{m-1}}, K_{m}=\frac{d_{m}-c_{m} K_{m-1}}{b_{m}+c_{m} L_{m-1}}$, is not accumulated (see [5]).
b. Discretizing the equation (4.4) by a difference scheme:

$$
\begin{aligned}
\Lambda_{2} C^{k+1}= & \frac{( \pm v+|v|)_{m, n}^{k+1}}{2} \frac{\left(C_{m, n}^{k+1}-C_{m, n-1}^{k+1}\right)}{\Delta y}+\frac{( \pm v-|v|)_{m, n}^{k+1}}{2} \frac{\left(C_{m, n+1}^{k+1}-C_{m, n}^{k+1}\right)}{\Delta y} \\
& -\gamma \frac{\left(C_{m, n+1}^{k+1}-2 C_{m, n}^{k+1}+C_{m, n-1}^{k+1}\right)}{\Delta y^{2}}+\frac{\sigma}{2} C_{m, n}^{k+1}
\end{aligned}
$$

we also get:

$$
\begin{equation*}
\tilde{a}_{n} C_{m, n+1}^{k+1}+\tilde{b}_{n} C_{m, n}^{k+1}+\tilde{c}_{n} C_{m, n-1}^{k+1}=\tilde{d}_{n} \tag{4,7}
\end{equation*}
$$

where,

$$
\begin{aligned}
& \tilde{a}_{m}=\frac{( \pm v-|v|)_{m, n}^{k+1} \theta d t}{2 \Delta y}-\frac{\gamma \theta d t}{(\Delta y)^{2}}, \quad \tilde{b}_{m}=1+\frac{\theta|v|_{m, n}^{k+1} d t}{\Delta y}+2 \frac{\gamma \theta d t}{(\Delta y)^{2}}+\frac{d t \sigma}{2} \\
& \tilde{c}_{m}=-\frac{( \pm v+|v|)_{m, n}^{k+1} d t}{2 \Delta y}-\frac{\gamma \theta d t}{(\Delta y)^{2}}, \quad \tilde{d}_{m}=C_{m, n}^{k+1 / 2}
\end{aligned}
$$

Obviously: $\tilde{b}_{m}>0, \tilde{a}_{m}<0, \tilde{c}_{m}<0$ and $\left|\tilde{b}_{m}\right| \geq\left|\tilde{a}_{m}\right|+\left|\tilde{c}_{m}\right|+\delta, \delta>0$.
Also, the equation system (4.7) has the unique solution and the double sweep method (4.6) does not produce an accumulated computational error.

## 5. Numerical experiments

The mentioned-above algorithm is applied to solve the following optimization problem of plant location:

- The computed rectangular region $G=1000 m \times 1000 m$ is covered by a uniform grid $51 \times 51$ with spacing steps: $d x=20 m, d y=20 m$.
- A constant velocity field $(u, v): u=0.5 \mathrm{~m} / \mathrm{s}, v=-0.5 \mathrm{~m} / \mathrm{s}$.
- Diffusion coefficient : $\gamma=0.5 \mathrm{~m}^{2} / \mathrm{s}$.
- Decay coefficient: $\sigma=0.0005 \mathrm{~s}^{-1}$.
- Time step: $d t=5 \mathrm{~s}$.
- Time simulation: $T=20000 \mathrm{~s}$.
- 3 considered sensitive rectangular areas $G_{k}$ inside $G(k=1,2,3)$ with the left-bottom corner coordinates and the right-top corner coordinates are as follows:
$+G_{1}=[(24.5,8.5),(25.5,9.5)],+G_{2}=[(37.5,12.5),(39.5,14.5)]$,
$+G_{3}=[(29.5,33.5),(30.5,34.5)]$.
- Standard concentration: $\overline{\bar{c}}_{k}=10 \mathrm{mg} / \mathrm{l}(k=1,2,3)$.

The numerical results are illustrated in Fig. 1. In this figure, the figure on the


Fig. 1. Distribution of value of the pollution level-reflecting functionals $Y_{k}^{*}$ for problem 1
contour lines indicates value of the pollution-level reflecting functionals $Y_{k}^{*}$. As a result, the domain $\Omega$ where the plant can be located so that the sanitary condition in the all areas $G_{k}$ are satisfied (that means $Y_{k}^{*} \leq \overline{\bar{c}}_{k}$ ) is in white.

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## BÀ̀ TOÁN 2 CHIỀU TỐI UUU XÁC ĐỊNH VI TRÍ NGUỒN THẢI

Bài báo trình bày các vấn đề sau: Bài toán liên hợp với bài toán lan truyền vật chất 2 chiều. Thuật toán xác định miền có thể đặt xí nghiệp sao cho phiếm hả̀m biểu thị mức độ ô nhiểm không vượt quá mức dộ cho phẹ́p ơ các vùng nhay cảm quan tầm. Đã áp dụng thuật toán này để tính toán cho một bài toán mẫu.

