

LEPTON FLAVOR VIOLATION IN ECONOMICAL 3-3-1 MODEL WITH NEUTRINO SINGLETS

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Abstract. *Lepton flavor violating (LFV) decays of charged leptons $\mu \rightarrow e\gamma$ and SM-like Higgs boson $h \rightarrow \mu\tau$ are discussed in the framework of the economical 3-3-1 model adding three additional singlets neutrinos ($E331X_R$). We will show that the appearance of new neutrinos and charged Higgs bosons in this model can result in interesting regions of the parameter space, which satisfy the recent experimental bound of $Br(\mu \rightarrow e\gamma)$ as well as give large $Br(h \rightarrow \mu\tau) \geq \mathcal{O}(10^{-7})$. In these regions, the charged Higgs boson mass lies within the range of 2 – 3 TeV. Our results show that the LFV decays are important channels to distinguish the $E331X_R$ and the original economical ($E331$) models, which predict suppressed LFV decay rates.*

Keywords: electroweak interactions, models beyond the Standard Model, non-standard-model neutrinos, right-handed neutrinos.

Classification numbers: 12.15.-y, 12.60.-i, 14.60.St.

I. INTRODUCTION

The economical 331 models [1, 2] are the simpler versions of the model with right-handed neutrino (331RN) [3], where only two Higgs triplets are introduced to generate masses of particles. These E331 models predict new neutrinos with very light masses, result in suppressed LFV processes like $\mu \rightarrow e\gamma$ and $h \rightarrow \mu\tau$, which are now being searched experimentally as signal of new physics beyond the Standard Model (SM) [4, 5]. Recently, new versions of the 331RN, where new singlet neutrinos are included, predicted heavy neutrinos through the inverse seesaw (ISS) mechanism [6, 7]. These models also predict that the $\text{Br}(\mu \rightarrow e\gamma)$ can reach the recent experimental bound, and the LFV decay of the SM-like Higgs boson $\text{Br}(h \rightarrow \mu\tau)$ may be large [7]. In this work, we will investigate the two mentioned LFV decays in the E331 X_R . Because of the simpler Higgs sectors than the previous 3-3-1 versions with ISS neutrinos, the LFV signal predicted by the E331 X_R can not be estimated based on the previous investigations. Therefore a detailed discussion in this work is very important and useful.

Our paper is arranged as follows. In Sec. II, we review ingredients of the E331 model relevant with the LFV processes. After that, we focus on important content of the E331 X_R , including ISS mechanism for generating neutrino masses, LFV couplings, and branching ratio of the two decays $\mu \rightarrow e\gamma$ and $h \rightarrow \mu\tau$. In Sec. III, numerical illustrations are presented. Section IV is our inclusion, where main results are given.

II. THE ECONOMICAL 3-3-1 MODEL WITH THREE NEUTRINO SINGLETS

II.1. The original E331 model

The economical 3-3-1 model (E331) was constructed based on the gauge group $SU(3)_C \times SU(3)_L \times U(1)_X$ with two scalar triplets [1, 2]. We review here important

$$\rho = \begin{pmatrix} \rho_1^+ \\ \rho_2^0 \\ \rho_3^- \end{pmatrix} \sim (3, \frac{2}{3}); \quad \phi = \begin{pmatrix} \phi_1^0 \\ \phi_2^- \\ \phi_3^0 \end{pmatrix} \sim (3, \frac{-1}{3}) \quad (1)$$

with vacuum expectation values (VEVs) are $\langle \phi_1^0 \rangle = u/\sqrt{2}$, $\langle \phi_3^0 \rangle = w/\sqrt{2}$, and $\langle \rho_2^0 \rangle = v/\sqrt{2} = 174$ GeV. It is enough to consider here $u = 0$, where the neutrinos masses will arise from the ISS mechanism.

Apart from photon, the E331 contains eight massive gauge bosons after spontaneous breaking, namely four neutral bosons (X^0, X^{0*}), (Z, Z') and four singly charged ones (W^\pm, V^\pm). The masses used in this work are $m_W = \frac{g_V}{2}$; $m_X = \frac{g_V}{2}$; $m_V = \frac{g\sqrt{v^2+w^2}}{2}$. The $SU(3)_L$ scale predicts new heavy particles, leading to the condition $w \gg v$.

The Higgs potential with only two scalar triplets is

$$\mathcal{V}_H = \mu_1^2 \phi^\dagger \phi + \mu_2^2 \rho^\dagger \rho + \lambda_1 (\phi^\dagger \phi)^2 + \lambda_2 (\rho^\dagger \rho)^2 + \lambda_3 (\phi^\dagger \phi) (\rho^\dagger \rho) + \lambda_4 (\rho^\dagger \phi) (\phi^\dagger \rho). \quad (2)$$

The E331 has a simple physical Higgs spectrum: two charged Higgs boson H_2^\pm , two CP-even Higgs (h, H) and one CP-odd boson A . Their masses and mixing are given in detailed in Ref [2], where the Higgs boson h is identified with that found by LHC. It is named the Standard Model-like (SM-like) Higgs boson because it has the similar property of the neutral Higgs boson predicted by

the SM. The masses of Higgs bosons relevant with this work are:

$$m_h^2 = v^2 \left[\left(2\lambda_2 - \frac{\lambda_3}{2\lambda_1} \right) + \mathcal{O} \left(\frac{v^2}{w^2} \right) \right], \quad m_{H_{2\pm}}^2 = \frac{\lambda_4}{g} (m_W^2 + m_X^2), \quad (3)$$

In the E331 model, neutral leptons are all included in triplets, $\psi_{aL} = (v_a, e_a, N_a)_L^T$, $a = 1, 2, 3$. The SM-like neutrino is v_a , and the new ones predicted by the E331 are N_{aL} . The lepton masses in the E331 model come from the Yukawa part

$$\mathcal{L}_l^Y = -h_{ab}^e \overline{\psi_{aL}} \rho e_{bR} + h_{ab}^v \epsilon^{ijk} \overline{(\psi_{aL})_i} (\psi_{bL})_j^c \rho_k^* + \text{h.c.}, \quad (4)$$

It was shown that the neutrinos in this model are all very light. Hence they result in very suppressed signals of LFV processes, which can not be detected by future experiments. In the next section, we will discuss the model containing heavy neutrinos.

In this work, we just focus on the decays relating with SM-like Higgs bosons and leptons, which are irrelevant with strong interactions like the quarks. The quark sector of the E331 model was discussed in detail previously [2], hence we will not repeat again. As we will see, the E331 X_R model is different from the E331 model by the lepton sector. Therefore, in the following it will be constructed thoroughly for determining the LFV couplings.

II.2. The E331 X_R model

In this work, the E331 X_R is constructed by adding three new singlets of right-handed neutrinos $X_{aR} \sim (1, 1, 0)$, $a=1,2,3$. The new Lagrangian part different from the E331 model is

$$-\mathcal{L}_{X_R} = Y_{ab} \overline{\psi_{aL}} \phi X_{bR} + \frac{1}{2} (\mu_X)_{ab} \overline{(X_{aR})^c} X_{bR} + \text{h.c.}. \quad (5)$$

Lagrangian (5) plays very important role to create masses and mixing of neutrinos following the ISS mechanism at tree level.

It can be seen that the E331 X_R has total nine neutrinos including three SM-like light v_a , three exotic neutrinos at the bottom of lepton triplets N_a and three right-handed neutrinos X_{aR} .

To determine neutrino masses and mixing, we will use the basis $v'_L = (v_L, N_L, (X_R)^c)^T$ and $(v'_L)^c = ((v_L)^c, (N_L)^c, X_R)^T$ with $v_R = ((v_{1L})^c, (v_{2L})^c, (v_{3L})^c)^T$, $N_R = ((N_{1L})^c, (N_{2L})^c, (N_{3L})^c)^T$, and $X_L = ((X_{1R})^c, (X_{2R})^c, (X_{3R})^c)^T$. Combining two Yukawa parts (4) and (5), the neutrino mass term is given as

$$-\mathcal{L}_{\text{mass}}^v = \frac{1}{2} \overline{v'_L} M^v (v'_L)^c + \text{h.c.}, \quad (6)$$

where M^v is the 9×9 matrix that is written in the block form as follows

$$M^v = \begin{pmatrix} 0 & m_D & 0 \\ m_D^T & 0 & M_R \\ 0 & M_R^T & \mu_X \end{pmatrix}. \quad (7)$$

In the expression (7), M_R is a 3×3 matrix having $(M_R)_{ab} \equiv Y_{ab} \frac{w}{\sqrt{2}}$ with $a, b = 1, 2, 3$, μ_X is also a 3×3 matrix was derived in (5), m_D is the Dirac mass matrix of neutrinos, which can be calculated in the same way shown in Ref. [7],

$$m_D = z \begin{pmatrix} 0 & x_{12} & x_{13} \\ -x_{12} & 0 & 1 \\ -x_{13} & -1 & 0 \end{pmatrix}. \quad (8)$$

Here, z , x_{12} , x_{13} are independent parameters. The mass matrix (7) has the form which creates neutrino mass from the ISS mechanism.

The mixing matrix of neutrino U^V satisfies the following condition

$$U^{VT} M^V U^V = \hat{M}^V = \text{diag}(\hat{m}_V, \hat{M}_N), \quad (9)$$

where $\hat{m}_V = \text{diag}(m_{n_1}, m_{n_2}, m_{n_3})$ and $\hat{M}_N = \text{diag}(m_{n_4}, m_{n_5}, \dots, m_{n_9})$, corresponding to the masses of active neutrinos and six extra neutrinos. The ISS mechanism guarantees that m_{n_a} are light and can be fitted with the current neutrino data. On the other hand, six exotic neutrino masses have order of $|M_R| \sim w$, hence they are very heavy.

II.3. Relevant LFV couplings and decays

The LFV coupling relating with gauge bosons can be traced from previous work,

$$\begin{aligned} \mathcal{L}^{\ell\ell V} &= \bar{\Psi}_{aL} \gamma^\mu D_\mu \Psi_{aL} \\ &= \frac{g}{\sqrt{2}} \left[U_{ai}^{V*} \bar{e}_a \gamma^\mu P_L n_i W_\mu^- + U_{ai}^V \bar{n}_i \gamma^\mu P_L e_a W_\mu^+ + U_{(a+3)i}^{V*} \bar{e}_a \gamma^\mu P_L n_i Y_\mu^- + U_{(a+3)i}^V \bar{n}_i \gamma^\mu P_L e_a Y_\mu^+ \right] \end{aligned}$$

The LFV couplings relating with Higgs bosons are changed because of the new Yukawa couplings (5). The Yukawa part (4) gives

$$\begin{aligned} &- h_{ab}^e \bar{\Psi}_{aL} \rho e_{bR} + h_{ab}^v \epsilon^{ijk} (\bar{\Psi}_{aL})_i (\Psi_{bL})_j \rho_k^* + \text{h.c.} \\ &\supset - \frac{g m_a}{\sqrt{2} m_W} \left[\frac{c_\xi}{\sqrt{2}} h \bar{e}_a e_a + c_\alpha \left(U_{(a+3)i}^V \bar{n}_i P_R e_a H_2^+ + U_{(a+3)i}^{V*} \bar{e}_a P_L n_i H_2^- \right) \right] \\ &- \frac{g c_\xi}{2 m_W} h \left[\sum_{c=1}^3 U_{ci}^V U_{cj}^{V*} \bar{n}_i (m_{n_i} P_L + m_{n_j} P_R) n_j \right] - \frac{g c_\alpha}{\sqrt{2} m_W} \left[U_{aj}^{V*} m_{n_j} \left(U_{j(b+3)}^{V\dagger} U_{bi}^V \right) H_2^- \bar{e}_a P_R n_i + \text{h.c.} \right], \end{aligned}$$

where two mixing parameters α and ξ are defined as

$$t_\alpha = \frac{m_W}{m_X} = \frac{v}{w}; \quad t_{2\xi} = \frac{\lambda_3 m_X m_W}{\lambda_1 m_X^2 - \lambda_2 m_W^2}. \quad (10)$$

In (5), the couplings from first term,

$$\begin{aligned} &- Y_{ab} \bar{\Psi}_{aL} \phi X_{bR} + \text{h.c.} \\ &\supset \frac{g}{2 m_W} (M_R)_{ab} \left[h s_\xi t_\alpha U_{(a+3)i}^V U_{(b+6)j}^V \bar{n}_i P_R n_j - \left(\sqrt{2} c_\alpha t_\alpha^2 U_{(b+6)i}^V \bar{e}_a P_R n_i H_2^- + \text{h.c.} \right) \right]. \quad (11) \end{aligned}$$

For convenience, we define

$$\begin{aligned} \lambda_{ij}^0 &= \sum_{c=1}^3 (U_{ci}^V U_{cj}^{V*} m_{n_i} + U_{ci}^{V*} U_{cj}^V m_{n_j}) - \sum_{c,d=1}^3 t_\alpha t_\xi (M_R^*)_{cd} \left[U_{(c+3)i}^{V*} U_{(d+6)j}^V + U_{(c+3)j}^{V*} U_{(d+6)i}^V \right], \\ \lambda_{ai}^R &= m_a U_{(a+3)i}^V, \quad \lambda_{ai}^L = \sum_{c=1}^3 (m_D^*)_{ac} U_{ci}^{V*} + t_\alpha^2 (M_R^*)_{ac} U_{(c+6)i}^{V*}. \end{aligned}$$

Based on the above calculation, all LFV couplings are listed in Table 1.

The partial decay width of the SM-like Higgs boson decay $h \rightarrow e_a e_b$ is [8, 9]

$$\Gamma(h \rightarrow e_a e_b) \equiv \Gamma(h \rightarrow e_a^- e_b^+) + \Gamma(h \rightarrow e_a^+ e_b^-) = \frac{m_h}{8\pi} (|\Delta_{(ab)L}|^2 + |\Delta_{(ab)R}|^2). \quad (12)$$

Table 1. Couplings relating with decays $h \rightarrow e_a e_b$ in the E331 X_R model. All momenta in the Feynman rules corresponding to these vertices are incoming.

Vertex	Factor
$h\bar{e}_a e_a$	$-\frac{igm_a}{2m_W} c_\xi$
$h\bar{n}_i n_j$	$-\frac{igc_\xi}{2m_W} \left(\lambda_{ij}^0 P_L + \lambda_{ij}^{0*} P_R \right)$
$H_2^- \bar{e}_a n_i, H_2^+ \bar{n}_i e_b$	$-\frac{igc_\alpha}{\sqrt{2}m_W} (\lambda_{ai}^{L*} P_R + \lambda_{ai}^{R*} P_L), -\frac{igc_\alpha}{\sqrt{2}m_W} (\lambda_{bi}^L P_L + \lambda_{bi}^R P_R)$
$W_\mu^+ e_b \bar{n}_i, W_\mu^- \bar{e}_a n_i$	$\frac{ig}{\sqrt{2}} U_{bi}^V \gamma^\mu P_L, \frac{ig}{\sqrt{2}} U_{ai}^{V*} \gamma^\mu P_L$
$V_\mu^+ \bar{n}_i e_b, V_\mu^- n_i \bar{e}_a$	$\frac{ig}{\sqrt{2}} U_{(b+3)i}^V \gamma^\mu P_L, \frac{ig}{\sqrt{2}} U_{(a+3)i}^{V*} \gamma^\mu P_L$
$hW^{+\mu} W^{-\nu}$	$igc_\xi m_W g_{\mu\nu}$
$hV^{+\mu} V^{-\nu}$	$ig m_V g_{\mu\nu} \sin(\alpha - \xi)$
$H_2^+ hV_\mu^-, V_\mu^+ H_2^- h$	$\frac{ig}{2} (P_{H_2^+} - P_h)_\mu \cos(\alpha - \xi), \frac{ig}{2} (P_h - P_{H_2^-})_\mu \cos(\alpha - \xi)$
$hH_2^+ H_2^-$	$i\lambda_{H_2}^{0\pm} = ic_\alpha^2 c_\xi [2\lambda_2 v - (\lambda_3 + \lambda_4) t_\xi w]$

The corresponding branching ratio is $\text{Br}(h \rightarrow e_a e_b) = \Gamma(h \rightarrow e_a e_b) / \Gamma_h^{\text{total}}$, where $\Gamma_h^{\text{total}} \simeq 4.1 \times 10^{-3}$ GeV is the total decay width [8, 10]. The $\Delta_{L,R}$ are scalar factors written as

$$\Delta_{L,R} = \sum_{i=1,5,7,8} \Delta_{L,R}^{(i)W} + \sum_{i=1}^{10} \Delta_{L,R}^{(i)Y}, \quad (13)$$

All private factors $\Delta_{L,R}^{(i)W}$ and $\Delta_{L,R}^{(i)Y}$ denote one-loop contributions from Feynman diagram in Fig. 1. Analytic expressions are computed using techniques presented in [11–13]. The results are listed in the appendix.

At one-loop level, the Feynman diagrams contributing to the decay $\mu \rightarrow e\gamma$ are shown in Fig. 2.

The formula of $\text{Br}(\mu \rightarrow e\gamma)$ is [14]

$$\text{Br}(\mu \rightarrow e\gamma) = \frac{12\pi^2}{G_F^2} |D_R|^2, \quad (14)$$

where $G_F = g^2 / (4\sqrt{2}m_W^2)$ and $D_R = D_R^W + D_R^V + D_R^{H_2^\pm}$,

$$\begin{aligned} D_R^W &= -\frac{eg^2}{32\pi^2 m_W^2} \sum_{j=1}^9 U_{1j}^{V*} U_{2j}^V F(t_{jW}), \quad D_R^V = -\frac{eg^2}{32\pi^2 m_V^2} \sum_{j=1}^9 U_{4j}^{V*} U_{5j}^V F(t_{jV}), \\ D_R^{H_2^\pm} &= -\frac{eg^2 c_\theta^2}{16\pi^2 m_W^2} \sum_{j=1}^9 \left[\frac{\lambda_{1j}^{L*} \lambda_{2j}^L}{m_{H_2^\pm}^2} \times \frac{1 - 6t_{jH} + 3t_{jH}^2 + 2t_{jH}^3 - 6t_{jH}^2 \ln(t_{jH})}{12(t_{jH} - 1)^4} \right. \\ &\quad \left. + \frac{m_{n_j} \lambda_{1j}^{L*} \lambda_{2j}^{R'}}{m_{H_2^\pm}^2} \times \frac{-1 + t_{jH}^2 - 2t_{jH} \ln(t_{jH})}{2(t_{jH} - 1)^3} \right], \end{aligned} \quad (15)$$

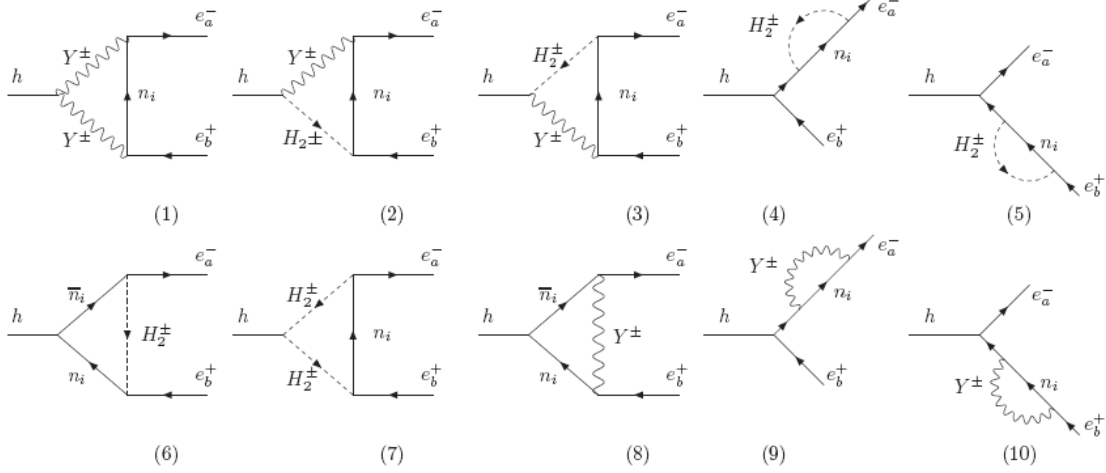


Fig. 1. One-loop Feynman diagrams contributing to the decay $h \rightarrow e_a e_b$ in the unitary gauge. Here $Y^\pm = (W^\pm, V^\pm)$.

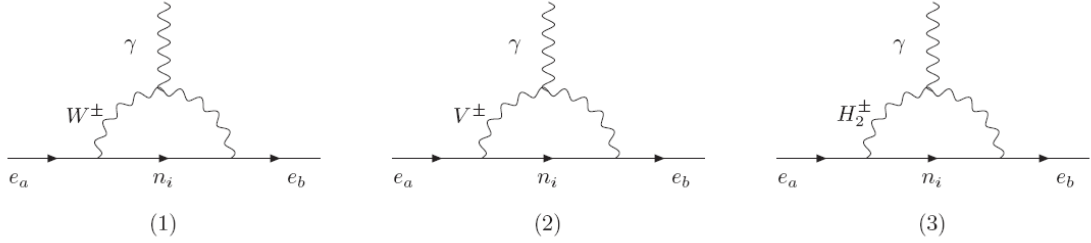


Fig. 2. Feynman diagrams for $\mu \rightarrow e\gamma$ at one-loop level in the unitary gauge.

where $t_{jW} \equiv \frac{m_{n_j}^2}{m_W^2}$, $t_{jV} \equiv \frac{m_{n_j}^2}{m_V^2}$, $t_{jH} \equiv \frac{m_{n_j}^2}{m_{H_2^\pm}^2}$, and

$$F(x) \equiv -\frac{10 - 43x + 78x^2 - 49x^3 + 4x^4 + 18x^3 \ln(x)}{12(x-1)^4}.$$

In the next section, we will focus on the allowed regions satisfying the upper bound of $\text{Br}(\mu \rightarrow e\gamma)$ and giving large $\text{Br}(h \rightarrow \mu\tau)$.

III. NUMERICAL RESULT

For numerical investigation, we choose a set of independent parameters as follows. The $SU(3)_L$ scale w is equivalent to the heavy charged boson mass $m_V = g\sqrt{v^2 + w^2}/2$; the singly charged Higgs mass $m_{H_2^\pm}$; the Higgs self-coupling factor λ_1 ; the mixing parameter ξ between

neutral Higgs bosons; heavy neutrino masses $m_{n_{4,5,6}}^2$; the parameters z of m_D matrix, and the matrix M_R .

Some parameters are the well-known experimental values [8]: the mass of W boson $m_W = 80.385$ GeV; charged lepton masses: $m_e = 5.10^{-4}$ GeV, $m_\mu = 0.105$ GeV, $m_\tau = 1.776$ GeV; $m_h = 125.1$ GeV; the $SU(2)_L$ gauge coupling $g = 0.651$, and $v \simeq 246$ GeV.

Other parameters are functions of the independent ones. Specifically, the mixing angle α and Higgs self couplings λ_4 are fixed by the relations (10) and (3), respectively. The Higgs couplings $\lambda_{2,3}$ are written as functions of ξ and m_h . The parameters $x_{1,2}$ and active neutrino masses in m_D are exactly the same as those mentioned in Ref. [7].

For simplicity and without losing generality, we choose $M_R = M_R I_3$ with $M_R = kz$, where k will be used. In numerical calculation, default values of free parameters are chosen as $\lambda_1 = 1$, $m_Y = 1.5$ TeV, $c_\xi = 0.99$, $\mu_X = 10^{-7}$ GeV. We note that the recent experimental data $h \rightarrow \gamma\gamma$ gives $|c_\xi| \geq 0.99$. The perturbative limit of Yukawa coupling of lepton gives $z \leq 600$ GeV. According to Ref. [7], in order to produce active neutrino masses and mixing which can be fitted to recent experimental data, we need a requirement that $k \geq 5.5$.

We checked numerically in the regions with large $k \geq 500$ and $z = 1, 5, 10, 500$ GeV, the results completely agree with [6, 7]. Furthermore, all allowed regions satisfying $Br(\mu \rightarrow e\gamma) \leq 4.2 \times 10^{-13}$ also give a small $Br(h \rightarrow \mu\tau) \leq 10^{-9}$. These regions are not interesting, we will not pay attention to them any more.

For illustration with small k , we will survey with two cases of $k = 6$ and $k = 9$ and different allowed values of $z = (50, 100, 300, 500, 600)$ GeV. As shown in figure 3, the results are very positive, where there exist narrow regions satisfying $Br(\mu \rightarrow e\gamma) \leq 4.2 \times 10^{-13}$. It can be explained

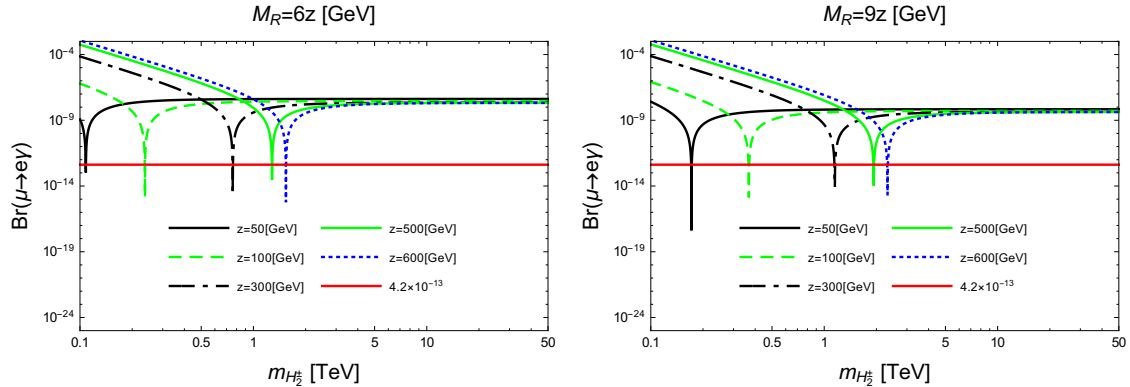


Fig. 3. $Br(\mu \rightarrow e\gamma)$ as function of $m_{H_2^\pm}$.

by the destructive correlations of the gauge and Higgs contributions. These properties does not appear in the ISS version of the SM, where charged Higgs bosons do not appear. The effects of k on the allowed regions can be seen in Fig. 4, where larger k will result in wider allowed regions.

To continue, we will investigate the $Br(h \rightarrow \mu\tau)$ in regions with small k and different values of z similarly to the case of $Br(\mu \rightarrow e\gamma)$. The results are shown in Fig. 5. The $Br(h \rightarrow \mu\tau)$ is enhanced strongly with increasing z , and changes slowly with $m_{H_2^\pm}$. $Br(h \rightarrow \mu\tau)$ can reach values

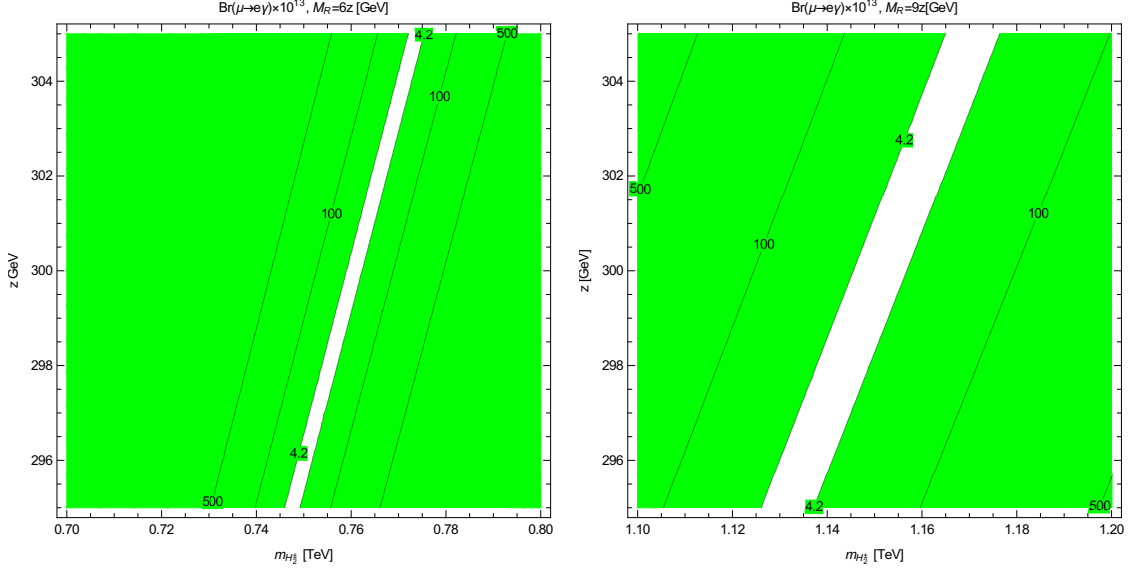


Fig. 4. Contours plots of $\text{Br}(\mu \rightarrow e\gamma)$ as functions of z and $m_{H_2^\pm}^2$. The green regions are excluded by the recent experimental data $\text{Br}(\mu \rightarrow e\gamma) \leq 4.2 \times 10^{-13}$.

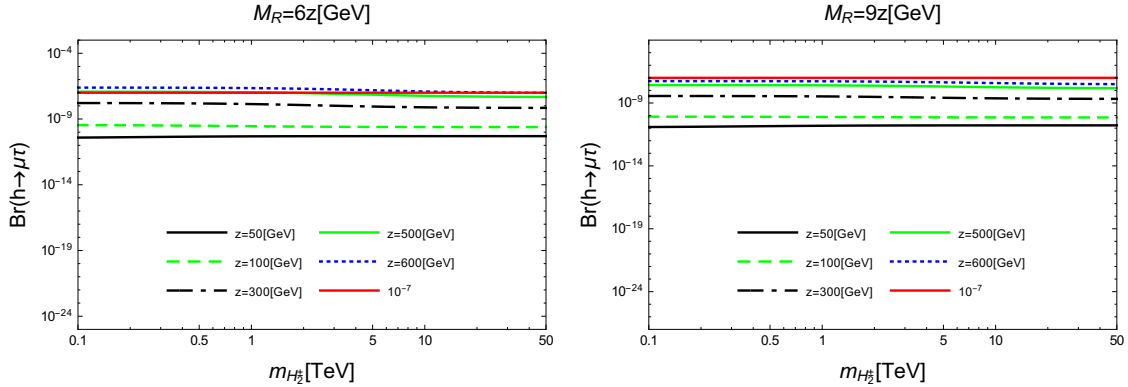


Fig. 5. Branching ratio of $h \rightarrow \mu\tau$ as function of $m_{H_2^\pm}$ with two case $k = 6$ (left panel) and $k = 9$ (right panel).

of $\mathcal{O}(10^{-7})$, in case $k = 6$ and $m_{H_2^\pm}$ is between $0.1 - 5$ TeV. With all values of z in the perturbative limit $z \leq 600$ GeV, there always exist allowed regions predicting small ranges of $m_{H_2^\pm}$.

The results of two specific cases $k = 6, k = 9$ and $z = 300; 500; 600$ GeV are given in Table 2. This is a very interesting prediction of the E331X_R, where both of the $\text{Br}(h \rightarrow \mu\tau)$ and $(\mu \rightarrow e\gamma)$ can reach the detection of the future experiments. These results are also consistent with previous publications such as [7, 14]. They are more interesting than those mentioned in [6], where only large $\text{Br}(\mu \rightarrow e\gamma)$ was shown in the case of large k .

Table 2. Values of $Br(h \rightarrow \mu \tau)$ with some fixed k and z .

k	z [GeV]	Value of $m_{H_2^\pm}$ [TeV] satisfying $Br(\mu \rightarrow e\gamma) \leq 4.2 \times 10^{-13}$	$Br(h \rightarrow \mu \tau)$
6	300	[0.759; 0.762]	$[1.41; 1.44] \times 10^{-8}$
9	300	[1.146; 1.156]	$[3.01; 3.30] \times 10^{-9}$
6	500	[1.280; 1.282]	$[1.024; 1.025] \times 10^{-7}$
9	500	[1.926; 1.933]	$[2.398; 2.400] \times 10^{-8}$
6	600	[1.540; 1.542]	$[2.1002; 2.1007] \times 10^{-7}$
9	600	[2.314; 2.323]	$[4.960; 4.963] \times 10^{-8}$

IV. CONCLUSIONS

In this paper we have shown that the E331 X_R predicts large Br of lepton flavor violation, especially in the regions having large m_D but light exotic neutrinos. Similarly to previous investigation of other 3-3-1 models with inverse seesaw neutrino, the region with very heavy exotic neutrino ($k \gg 1$), the model predicts that $Br(\mu \rightarrow e\gamma)$ can reach recent upper bound of experiment. In contrast, the $Br(h \rightarrow \mu \tau) < \mathcal{O}(10^{-9})$, which is still too small for searching this signal. For small $k \geq 6$, there exist very narrow regions satisfying experimental bound of $Br(\mu \rightarrow e\gamma)$ and predict large $Br(h \rightarrow \mu \tau)$, which can be larger than $\mathcal{O}(10^{-7})$. Furthermore, $m_{H_2^\pm}$ is in order of $\mathcal{O}(1)$ TeV. In the region with small k , the appearance of H_2^\pm is very important. Because the allowed regions, which have very small $Br(\mu \rightarrow e\gamma)$, arise from the destructive correlations of charged Higgs and gauge boson contributions to the $Br(\mu \rightarrow e\gamma)$.

The LFV decay rates of the two processes $\mu \rightarrow e\gamma$ and $h \rightarrow \mu \tau$ are suppressed in the E331 models. Hence, these two channel decays will give important information to distinguish the E331 models and the E331 X_R if any LFV signals are detected in the future.

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APPENDIX

Here we only give analytic expressions for contribution from Feynman diagrams (4), (5), (6), (7) in Fig.1 with charged Higgs bosons (H_2^\pm). The remaining diagrams are calculated in the same way.

$$\begin{aligned}
\Delta_L^{(4+5)VH_2^\pm} &= -\frac{g^3 c_\alpha c_\zeta^2}{32\pi^2 m_W^3 (m_a^2 - m_b^2)} \\
&\times \sum_{i=1}^9 \left[m_a m_b m_{n_i} \lambda_{ai}^{L*} \lambda_{bi}^R (B_0^{(1)} - B_0^{(2)}) + m_{n_i} \lambda_{ai}^{R*} \lambda_{bi}^L (m_b^2 B_0^{(1)} - m_a^2 B_0^{(2)}) \right. \\
&\quad \left. + m_a m_b (\lambda_{ai}^{L*} \lambda_{bi}^L m_b + \lambda_{ai}^{R*} \lambda_{bi}^R m_a) (B_1^{(1)} + B_1^{(2)}) \right], \\
\Delta_R^{(4+5)VH_2^\pm} &= -\frac{g^3 c_\alpha c_\zeta^2}{32\pi^2 m_W^3 (m_a^2 - m_b^2)} \\
&\times \sum_{i=1}^9 \left[m_a m_b m_{n_i} \lambda_{ai}^{R*} \lambda_{bi}^L (B_0^{(1)} - B_0^{(2)}) + m_{n_i} \lambda_{ai}^{L*} \lambda_{bi}^R (m_b^2 B_0^{(1)} - m_a^2 B_0^{(2)}) \right. \\
&\quad \left. + m_a m_b (\lambda_{ai}^{R*} \lambda_{bi}^R m_b + \lambda_{ai}^{L*} \lambda_{bi}^L m_a) (B_1^{(1)} + B_1^{(2)}) \right], \\
\Delta_L^{(6)VH_2^\pm} &= -\frac{g^3 c_\alpha c_\zeta^2}{32\pi^2 m_W^3} \sum_{i,j=1}^9 \left\{ \lambda_{ij}^{0*} \left[\lambda_{ai}^{R*} \lambda_{bj}^L (B_0^{(12)} + m_{H_k^\pm}^2 C_0 - m_a^2 C_1 + m_b^2 C_2) \right. \right. \\
&\quad \left. \left. + \lambda_{ai}^{R*} \lambda_{bj}^R m_b m_{n_j} C_2 - \lambda_{ai}^{L*} \lambda_{bj}^L m_a m_{n_i} C_1 \right] \right. \\
&\quad \left. + \lambda_{ij}^0 \left[\lambda_{ai}^{R*} \lambda_{bj}^L m_{n_i} m_{n_j} C_0 + \lambda_{ai}^{R*} \lambda_{bj}^R m_{n_i} m_b (C_0 + C_2) \right. \right. \\
&\quad \left. \left. + \lambda_{ai}^{L*} \lambda_{bj}^L m_a m_{n_j} (C_0 - C_1) + \lambda_{ai}^{L*} \lambda_{bj}^R m_a m_b (C_0 - C_1 + C_2) \right] \right\}, \\
\Delta_R^{(6)VH_2^\pm} &= -\frac{g^3 c_\alpha c_\zeta^2}{32\pi^2 m_W^3} \sum_{i,j=1}^9 \left\{ \lambda_{ij}^0 \left[\lambda_{ai}^{L*} \lambda_{bj}^R (B_0^{(12)} + m_{H_2^\pm}^2 C_0 - m_a^2 C_1 + m_b^2 C_2) \right. \right. \\
&\quad \left. \left. + \lambda_{ai}^{L*} \lambda_{bj}^L m_b m_{n_j} C_2 - \lambda_{ai}^{R*} \lambda_{bj}^R m_a m_{n_i} C_1 \right] \right. \\
&\quad \left. + \lambda_{ij}^{0*} \left[\lambda_{ai}^{L*} \lambda_{bj}^R m_{n_i} m_{n_j} C_0 + \lambda_{ai}^{L*} \lambda_{bj}^L m_{n_i} m_b (C_0 + C_2) \right. \right. \\
&\quad \left. \left. + \lambda_{ai}^{R*} \lambda_{bj}^R m_a m_{n_j} (C_0 - C_1) + \lambda_{ai}^{R*} \lambda_{bj}^L m_a m_b (C_0 - C_1 + C_2) \right] \right\}, \\
\Delta_L^{(7)VH_2^\pm} &= \frac{g^2 \lambda_{H_k^\pm} c_\zeta^2}{16\pi^2 m_W^2} \sum_{i=1}^9 \left[-\lambda_{ai}^{R*} \lambda_{bi}^L m_{n_i} C_0 - \lambda_{ai}^{L*} \lambda_{bi}^L m_a C_1 + \lambda_{ai}^{R*} \lambda_{bi}^R m_b C_2 \right], \\
\Delta_R^{(7)VH_2^\pm} &= \frac{g^2 \lambda_{H_2^\pm} c_\zeta^2}{16\pi^2 m_W^2} \sum_{i=1}^9 \left[-\lambda_{ai}^{L*} \lambda_{bi}^R m_{n_i} C_0 - \lambda_{ai}^{R*} \lambda_{bi}^R m_a C_1 + \lambda_{ai}^{L*} \lambda_{bi}^L m_b C_2 \right],
\end{aligned}$$